

fhn_snlc

December 18, 2015

0.1 Saddle-nodes of periodic orbits in the FitzHugh-Nagumo model

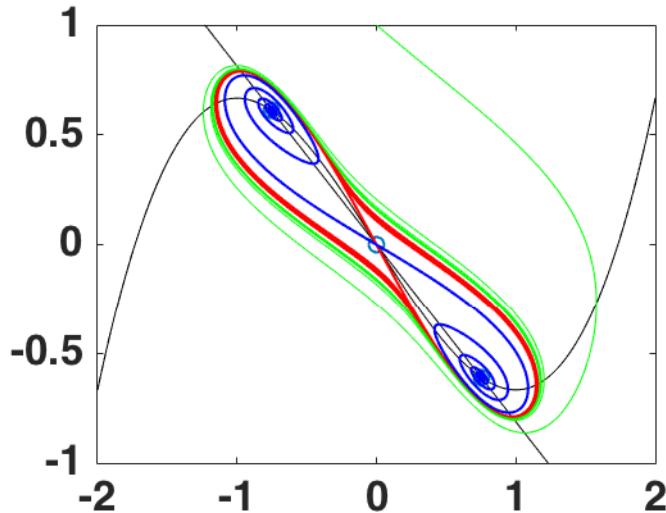
The version of the FitzHugh-Nagumo model used in this notebook is given by the equations

$$\dot{x} = y - x^3/3 + x\dot{y} = -e(ax + b + cy)$$

Newton's method is used on a return map and its derivative to locate a saddle-node of periodic orbits.

Plot phase portraits for parameters that straddle the bifurcations. Active parameter is e .

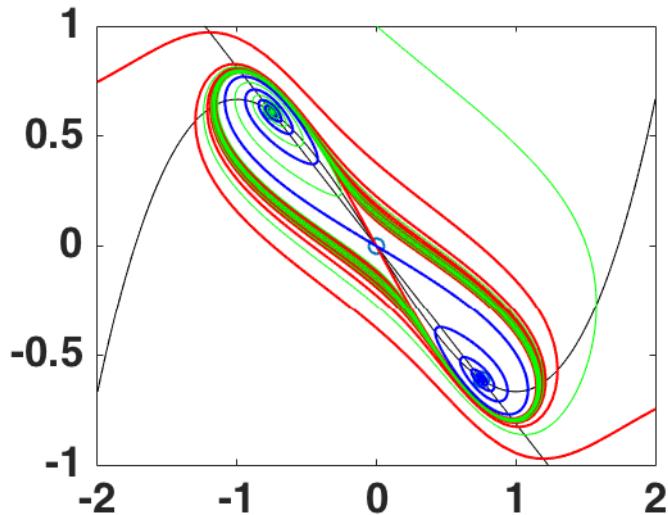
```
In [1]: p = [13/16; 0; 1; 0.544];
eqm = fhn_eq(p)
fhn_nullcline(p);
%Stable and unstable manifolds
[ws1,ws2,wu1,wu2] = fhn2_smflds(eqm(1,:),p,1e-5);
options = dopset('AbsTol',1e-14,'RelTol',1e-12,'MaxStepSize',0.01,'MaxSteps',1e7,'EventTol',1e-12);
[ts,pts,te,ye,ie,stats] = dop853('fhn2',[0 2000],[0,1],options,p);
plot(pts(:,1),pts(:,2),'g')
axis([-2,2,-1,1])
%
```



```
Out[1]: eqm =
```

$$\begin{matrix} 0 & 0 \\ 0.7500 & -0.6094 \\ -0.7500 & 0.6094 \end{matrix}$$

```
In [2]: p = [13/16; 0; 1; 0.545];
eqm = fhn_eq(p)
fhn_nullcline(p);
%Stable and unstable manifolds
[ws1,ws2,wu1,wu2] = fhn2_smflds(eqm(1,:),p,1e-5);
options = dopset('AbsTol',1e-14,'RelTol',1e-12,'MaxStepSize',0.01,'MaxSteps',1e7,'EventTol',1e-12);
[ts,pts,te,ye,ie,stats] = dop853('fhn2',[0 2000],[0,1],options,p);
plot(pts(:,1),pts(:,2),'g')
axis([-2,2,-1,1])
```



Out [2]: eqm =

$$\begin{matrix} 0 & 0 \\ 0.7500 & -0.6094 \\ -0.7500 & 0.6094 \end{matrix}$$

Exit of dop853 at $x = -2.8247865932943455e+02$, step size too small $h = -6.0144644841320101e-13$
 Exit of dop853 at $x = -2.8247865932943455e+02$, step size too small $h = -6.0144644841320101e-13$

Use bisection to get close to bifurcation. Then use return map to examine details.

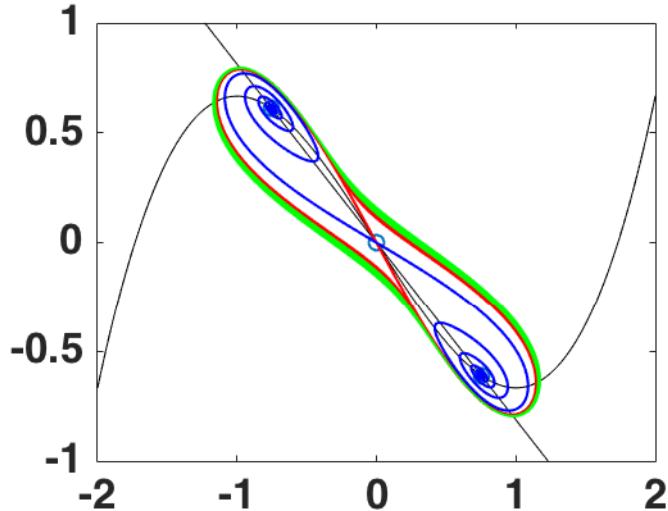
```
In [3]: %The parameters
p = [13/16; 0; 1; 0.5443639468]
axis([-2,2,-1,1])
%Equilibrium points and nullclines
eqm = fhn_eq(p)
fhn_nullcline(p);
%Stable and unstable manifolds
[ws1,ws2,wu1,wu2] = fhn2_smflds(eqm(1,:),p,1e-5);
axis([-2,2,-1,1])

%Plot 100 trajectories to first return with increasing coordinate for
%cross-section that determines linear nullcline
xn = 100;
```

```

%x coordinate increment between intial points
xinc = 8e-5;
%matrix of initial points
xin = -0.984 +8e-5*[1:100];
yin = (-p(1)*xin-p(2))./p(3);
pin = [xin',yin'];
%For output - checking that return is to correct region
pin2= [];
pout = [];
options = dopset('AbsTol',1e-14,'RelTol',1e-12,'MaxStepSize',0.01,'MaxSteps',1e7,'EventTol',1e-12);
%Compute the trajectoies and store initial and final points
for j=1:xn
    [ts,pts,te,ye,ie,stats] = dop853('fhn_el2',[0 2000],pin(j,:),options,p);
    plot(pts(:,1),pts(:,2), 'g')
    if (pts(end,1) <-0.75)
        pin2 = [pin2;pin(j,:)];
        pout = [pout;pts(end,:)];
    end
end

```



Out[3]: *p* =

```

0.8125
0
1.0000
0.5444

```

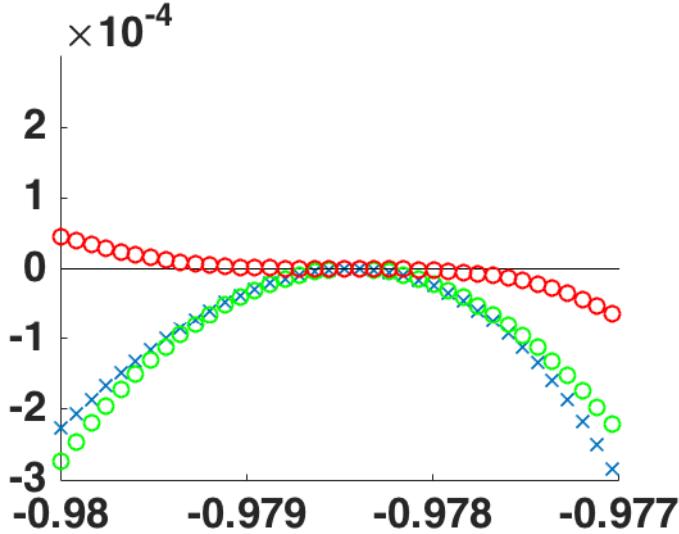
```

eqm =
0 0
0.7500 -0.6094
-0.7500 0.6094

```

Analyze the return data

```
In [4]: %Compute derivatives for return
%difference of x-coordinate endpoints
xd = pin2(:,1) -pout(:,1);
%Plot x-coordinate differences vs. initial values of x
figure(2)
clf
hold on
plot(pin2(:,1),xd,'x')
plot([pin2(1,1),pin2(end,1)],[0,0],'k')
%CLOSEST RETURN TO A FIXED POINT
[resid,ind] = min(abs(pin2-pout))
ind = ind(1);
%Compute first and second derivatives of x-coordinate endpoint differences
x0 = pin2(ind,1);
y0 = xd(ind,1)
dx0 = (xd(ind+1)-xd(ind-1))/(2*xinc)
dxx0 = (xd(ind+1)-2*xd(ind)+xd(ind-1))/(xinc^2)
%Quadratic fit to return
qfit = y0 + dx0*(pin2(:,1)-x0) + dxx0*(pin2(:,1)-x0).^2/2;
%Plot quadratic green
plot(pin2(:,1),qfit,'go')
%Plot residual red
plot(pin2(:,1),xd-qfit,'ro')
axis([-0.98,-0.977,-3e-4,3e-4])
```



```
Out[4]: resid =
1.0e-10 *
0.3646    0.2962

ind =
```

69 69

y0 =

3.6458e-11

dx0 =

0.0085

dxx0 =

-225.5491

Repeat for parameters without a periodic orbit

In [5]: %The parameters

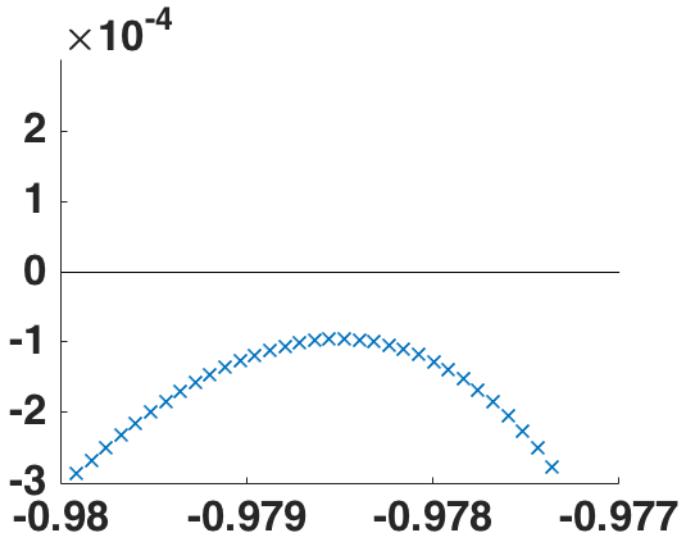
```
p = [13/16; 0; 1; 0.5444]
axis([-2,2,-1,1])
%Equilibrium points and nullclines
eqm = fhn_eq(p)
fhn_nullcline(p);
%Stable and unstable manifolds
[ws1,ws2,wu1,wu2] = fhn2_smflds(eqm(1,:),p,1e-5);
axis([-2,2,-1,1])

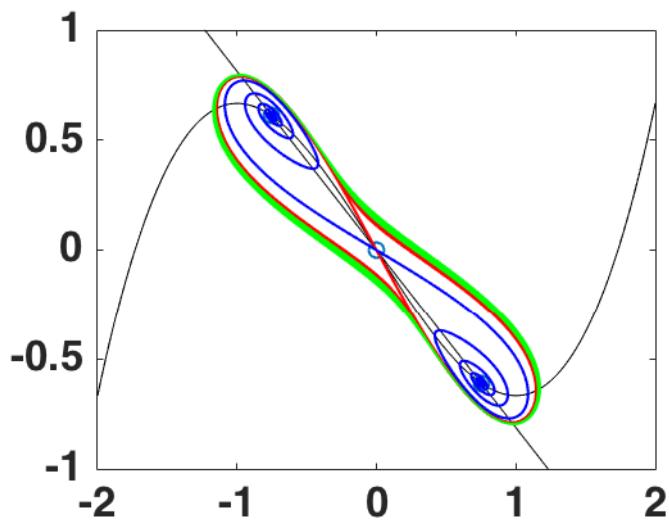
%Plot 100 trajectories to first return with increasing coordinate for
%cross-section that determines linear nullcline
xn = 100;
%x coordinate increment between intial points
xinc = 8e-5;
%matrix of initial points
xin = -0.984 +8e-5*[1:100];
yin = (-p(1)*xin-p(2))./p(3);
pin = [xin',yin'];
%For output - checking that return is to correct region
pin2= [];
pout = [];
options = dopset('AbsTol',1e-14,'RelTol',1e-12,'MaxStepSize',0.01,'MaxSteps',1e7,'EventTol',1e-12);
%Compute the trajectories and store initial and final points
for j=1:xn
    [ts,pts,te,ye,ie,stats] = dop853('fhn_el2',[0 2000],pin(j,:),options,p);
    plot(pts(:,1),pts(:,2),'g')
    if (pts(end,1) <-0.75)
        pin2 = [pin2;pin(j,:)];
        pout = [pout;pts(end,:)];
    end
end
%Compute derivatives for return
%difference of x-coordinate endpoints
xd = pin2(:,1)-pout(:,1);
%Plot x-coordinate differences vs. initial values of x
```

```

figure(2)
clf
hold on
plot(pin2(:,1),xd,'x')
plot([pin2(1,1),pin2(end,1)],[0,0],'k')
%Closes return to a fixed point
[resid,ind] = min(abs(pin2-pout))
ind = ind(1);
%Compute first and second derivatives of x-coordinate endpoint differences
x0 = pin2(ind,1);
y0 = xd(ind,1)
dx0 = (xd(ind+1)-xd(ind-1))/(2*xinc)
dxx0 = (xd(ind+1)-2*xd(ind)+xd(ind-1))/(xinc^2)
%Quadratic fit to return
qfit = y0 + dx0*(pin2(:,1)-x0) + dxx0*(pin2(:,1)-x0).^2/2;
%Plot quadratic green
%plot(pin2(:,1),qfit,'go')
%Plot residual red
%plot(pin2(:,1),xd-qfit,'ro')
axis([-0.98,-0.977,-3e-4,3e-4])

```





Out[5]: p =

```
0.8125
0
1.0000
0.5444
```

eqm =

```
0 0
0.7500 -0.6094
-0.7500 0.6094
```

resid =

```
1.0e-04 *
0.9654 0.7844
```

ind =

```
69 69
```

y0 =

```
-9.6541e-05
```

dx0 =

```
-0.0057
```

```
dxx0 =  
-231.3807
```

0.1.1 Newton iteration

Now we want to use Newton iteration to obtain a more accurate value of the parameter at the saddle-node of limit cycles. The linear nullcline is chosen as a cross-section. (Any periodic orbit must cross both nullclines.) Three initial points are chosen on the nullcline, separated by distances proportional to x_{inc} , defined above. The value of the return map σ at these points is computed, together with a centered difference estimate of σ' at the middle point. The defining equations for a saddle-node of limit cycles are that $\sigma(x) = x$ and $\sigma'(x) = 1$

In [6]: % Starting parameters and initial point

```
format long  
p0 = [13/16; 0; 1; 0.5443639468];  
z0 = [-0.97848,0.795015];  
p = p0  
z = z0  
% Finite difference increments  
xinc = 3e-4;  
pinc = 1e-5;  
% maximum number of Newton iterations  
nsteps = 20;  
resid = [];  
for j=1:nsteps  
    % Trajectories for computing x derivatives  
    options = dopset('AbsTol',1e-14,'RelTol',1e-12,'MaxStepSize',0.01,'MaxSteps',1e7,'EventTol',1e-12);  
    [ts,pts,te,ye,ie,stats] = dop853('fhn_el2',[0 2000],z,options,p);  
    zout = pts(end,:);  
    % Residual of sigma(x) - x  
    xd = zout(1)-z(1);  
    [ts,pts,te,ye,ie,stats] = dop853('fhn_el2',[0 2000],z+[xinc,-p(1)*xinc/p(3)],options,p);  
    zr = pts(end,:);  
    xr = zr(1) - z(1) - xinc;  
    [ts,pts,te,ye,ie,stats] = dop853('fhn_el2',[0 2000],z-[xinc,-p(1)*xinc/p(3)],options,p);  
    zl = pts(end,:);  
    xl = zl(1) - z(1) + xinc;  
    % Residual of sigma'(x) - 1  
    dx = (xr-xl)./(2*xinc);  
    % Accumulate residuals  
    resid = [resid;[xd,dx]];  
    % Convergence test  
    if abs(xd) < 1e-10 && abs(dx) < 1e-10  
        j=j  
        resid  
        z  
        p  
        return  
    end  
    % sigma''(x) for Jacobian  
    dxx = (zr(1)-2*zout(1)+zl(1))/(xinc^2);
```

```

%
% Increment parameter p(4)
pp = p+pinc*[0;0;0;1];
% Compute trajectories for parameters pp
[ts,pts,te,ye,ie,stats] = dop853('fhn_el2',[0 2000],z,options,pp);
zoutp = pts(end,:);
xdp = zoutp(1)-z(1);
[ts,pts,te,ye,ie,stats] = dop853('fhn_el2',[0 2000],z+[xinc,-p(1)*xinc/p(3)],options,pp);
zrp = pts(end,:);
xrp = zrp(1) - z(1) - xinc;
[ts,pts,te,ye,ie,stats] = dop853('fhn_el2',[0 2000],z-[xinc,-p(1)*xinc/p(3)],options,pp);
zlp = pts(end,:);
xlp = zlp(1) - z(1) + xinc;
dxp = (xrp-xlp)./(2*xinc);
% Derivatives with respect to p(4)
dpar = (xdp-xd)/pinc;
dxpar = (dxp-dx)/pinc;
%Jacobian for defining function (\sigma - id, \sigma' - 1)
jac = [[dx,dpar];[dxx,dxpar]];
%
% Newton step
newt_delta = jac\ [xd;dx];
xpnew = [z(1);p(4)] - newt_delta;
monitor = [xd,dx,newt_delta(2)];
% Update z and p
z = [xpnew(1),(-p(1)*xpnew(1)-p(2))/p(3)];
p = [p(1:3);xpnew(2)];
end
p0(4)-p(4)

```

Out[6]: p =

```

0.8125000000000000
0
1.0000000000000000
0.5443639468000000

```

z =

```
-0.9784800000000000 0.7950150000000000
```

monitor =

```
-0.000000000036458 -0.007128927614822 -0.000000083933687
```

monitor =

```
1.0e-04 *
```

```
0.000662913762772 0.638614027224970 0.000247311646362
```

```

monitor =
1.0e-07 *
0.004601613534660 -0.740449543765292 0.001717175267417

monitor =
1.0e-09 *
-0.000253908005732 0.539458343447319 -0.000094750367979

j =
5

resid =
-0.00000000036458 -0.007128927614822
0.000000066291376 0.000063861402722
0.000000000460161 -0.000000074044954
-0.000000000000254 0.000000000539458
0.000000000000008 -0.000000000005476

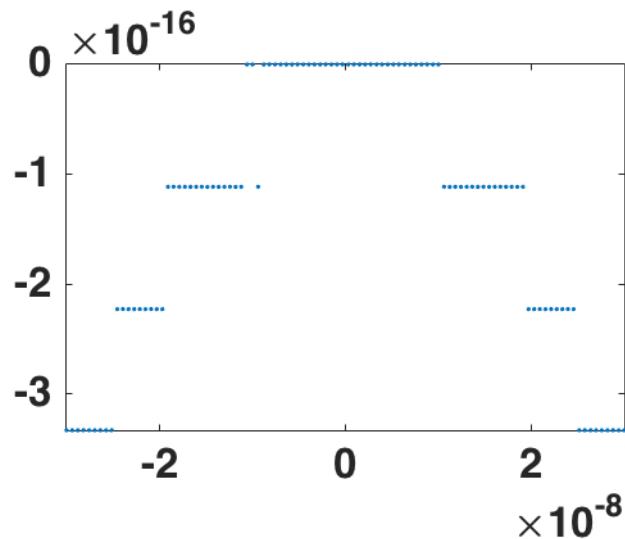
z =
-0.978448834514833 0.794989678043302

p =
0.812500000000000
0
1.000000000000000
0.544364005830899

```

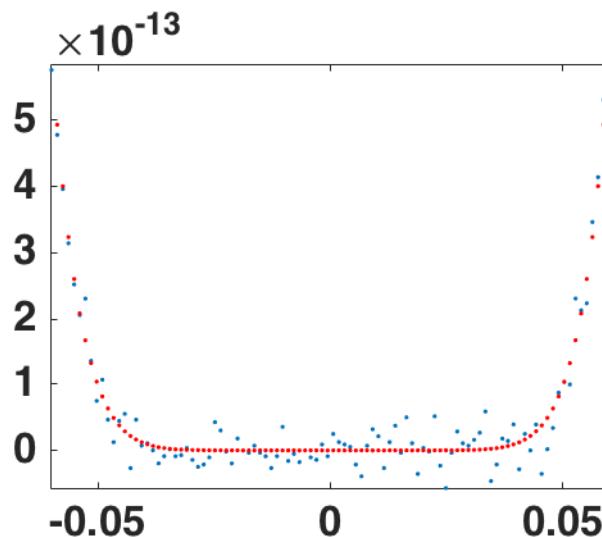
One problem with this Newton iteration is that numerical identification of *where* a function has an extreme value is difficult. Here are two examples that illustrate this difficulty. In the first, we plot computed values of the function $f(x) = \sin(x) - x/2$ and see that the discretization of floating point values makes the function appear constant over an interval of length comparable to 1e-8.

```
In [7]: xinc = 3e-8;
x = linspace(pi/3-xinc,pi/3+xinc);
y = sin(x) - x/2 - sqrt(3)/2+pi/6;
figure(1)
plot(x-pi/3,y,'.')
axis([-xinc,xinc,min(y),max(y)])
```



The second example writes $f(u) = (1-u)^{10}$ in its expanded form. The plot shows that cancellation in the round-off evaluation of different terms can produce “noisy” values while evaluation of the function in the unexpanded form does not.

```
In [8]: uinc = 0.06;
u = linspace(1-uinc,1+uinc);
v = u .^ 10 - 10 * u .^ 9 + 45 * u .^ 8 - 120 * u .^ 7 + 210 * u .^ 6 - 252 * u .^ 5 + 210 * u
w = (1-u).^10;
figure(2)
plot(u-1,v,'.',u-1,w,'r.')
axis([-uinc,uinc,min(v),max(v)])
min(v)
```



```
Out[8]: ans =
```

```
-5.684341886080801e-14
```

These examples suggest that it is hardly possible to locate precisely where the return map $\sigma(z)$ of the FitzHugh-Nagumo model has an extreme point for $\sigma(z) - z$. However, it does seem feasible to estimate accurately how the extreme value depends upon the parameter e . With this in mind, we leave at as a challenge to implemnt such an algorithm and embed it into a continuation scheme.

In []: