

Continuation and bifurcation analysis in experiments II

Jan Sieber

University of Exeter (UK)

Plan

- ▶ how to continue unstable branches of equilibria and periodic orbits in experiments

Lecture 1

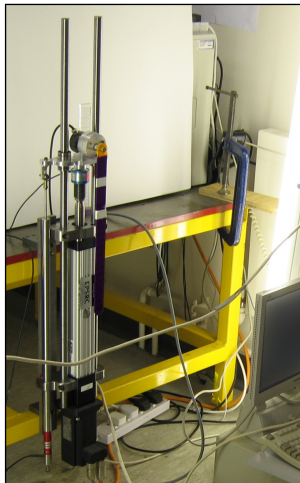
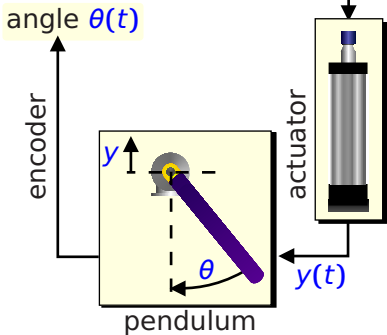
- ▶ motivating examples
 - ▶ short intro to concepts for feedback control
- ▶ ⇒ **Lecture 2** detailed examples (simple mechanical experiments)
 - + periodic orbits close to Hopf bifurcation
- ▶ **Lecture 3** Continuation for delay-differential equations (mostly done by Tony): looking under the hood
- ▶ **Afternoons & hands-on workshop** help with DDE-Biftool (for delay equations), AUTO, coco
- ▶ (for myself) learn more coco

Example: rotation in pendulum

Schema

excitation $p \sin(\omega t)$

angle $\theta(t)$



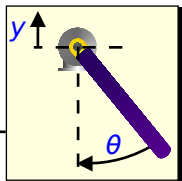
Example: rotation in pendulum

Schema

excitation $p \sin(\omega t)$

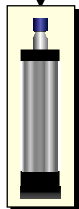
angle $\theta(t)$

encoder



pendulum

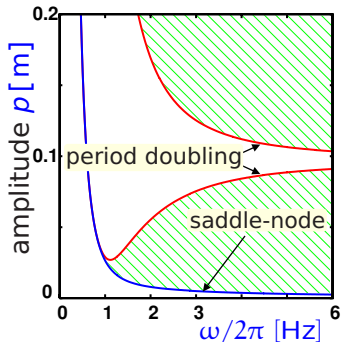
actuator



$y(t)$

Numerics

rotation

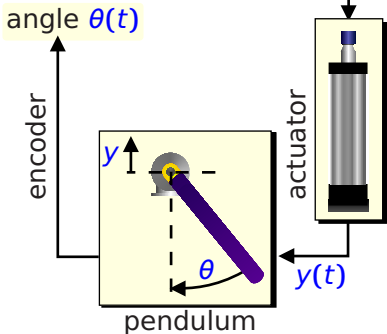


Example: rotation in pendulum

Schema

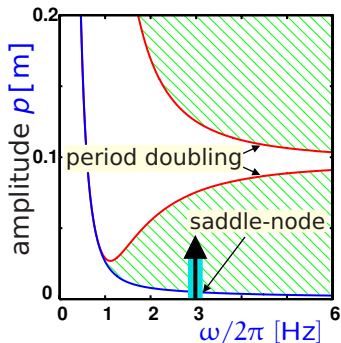
excitation $p \sin(\omega t)$

angle $\theta(t)$



Numerics

rotation

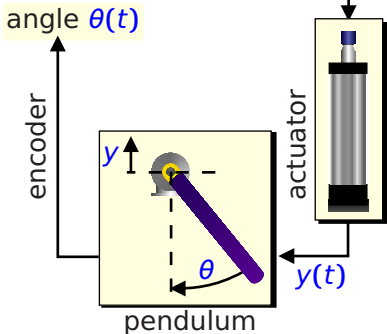


Example: rotation in pendulum

Schema

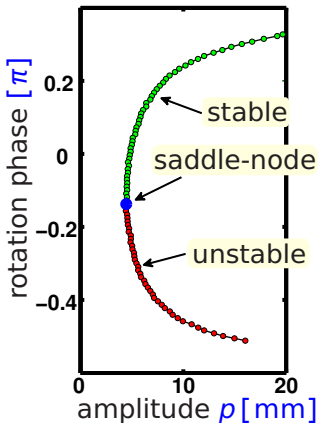
excitation $p \sin(\omega t)$

angle $\theta(t)$



Experiment

rotation

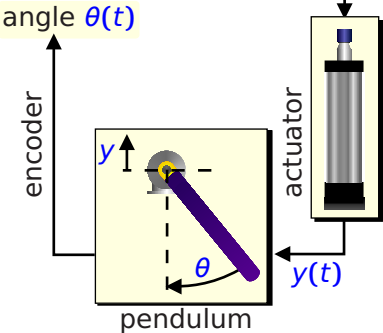


Example: rotation in pendulum

Schema

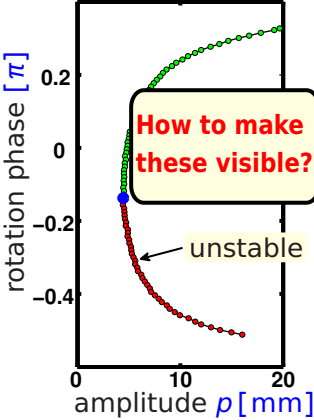
excitation $p \sin(\omega t)$

angle $\theta(t)$



Experiment

rotation



Example: rotation in pendulum

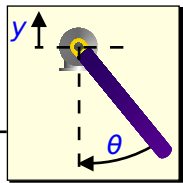
Schema

excitation $p \sin(\omega t)$

demand $\theta_d(t)$

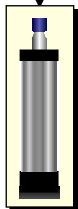
angle $\theta(t)$

encoder



pendulum

actuator



$y(t)$

-

PD

+

Control

- feed back $\theta - \theta_d$
- ⇒ always stable
- ⇒ output not natural
- How to choose θ_d ?
- ⇒ find θ_d such that $\theta_d - \theta = 0$

Computational side

Unknowns: θ_d (Fourier modes),
 p (parameter)

“Equation”: $\theta_d - \theta[\theta_d, p] = 0$

⇒ **use Newton iteration**

Newton iteration requires function $\text{eqs}(\theta_d, p)$:

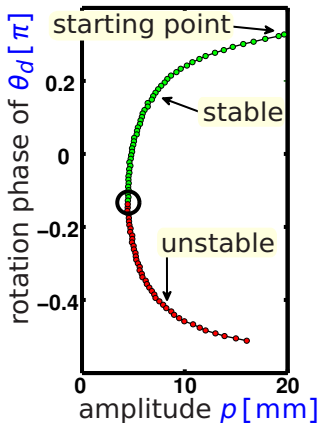
- ▶ set control demand to θ_d , excitation amplitude p
- ▶ wait until output θ is periodic again
- ▶ return Fourier modes of $\theta - \theta_d$

+ **pseudo-arclength condition**

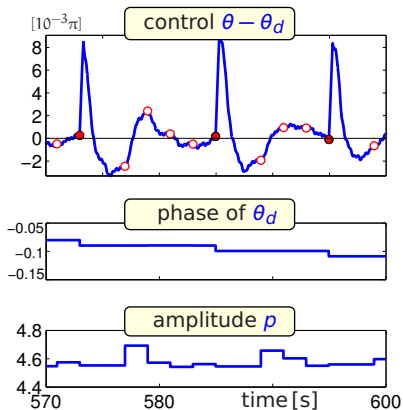
Experimental results

Experiment

bifurcation diagram



time profile of continuation



Experimental results

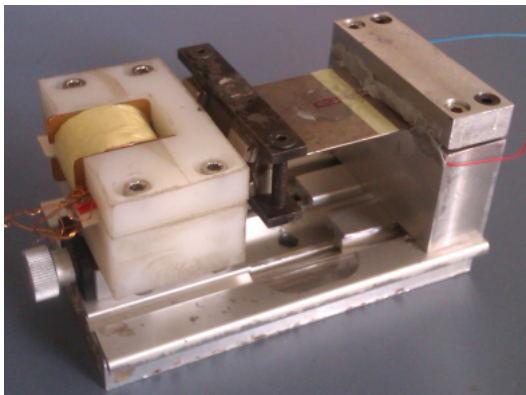
video

Comments

- ▶ model accurate except for damping, but damping has large influence on bifurcation diagram
- ▶ true dimension of nonlinear system for Newton iteration: 1 (phase shift ϕ)
other harmonics found with “time-delayed feedback”/Picard iteration

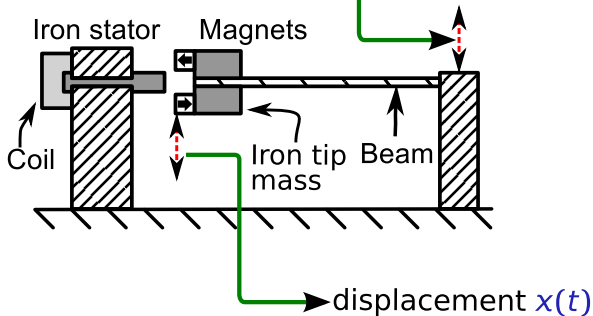
David Barton's results (Bristol)

Nonlinear Energy Harvester

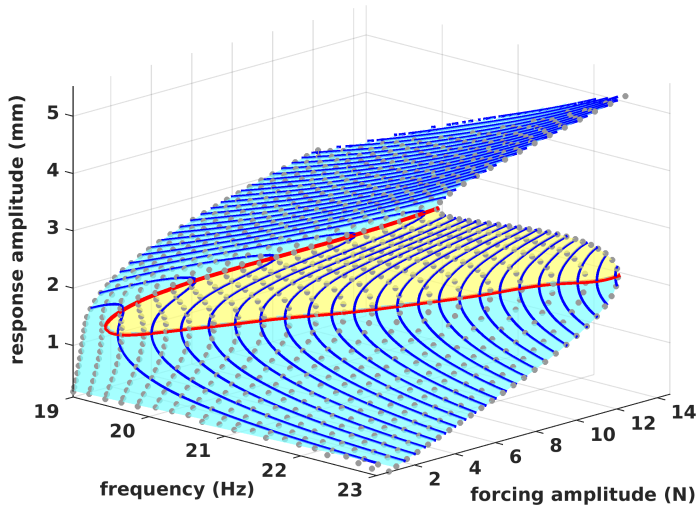


David Barton's results (Bristol)

excitation force $a \sin(\omega t) + u(t)$

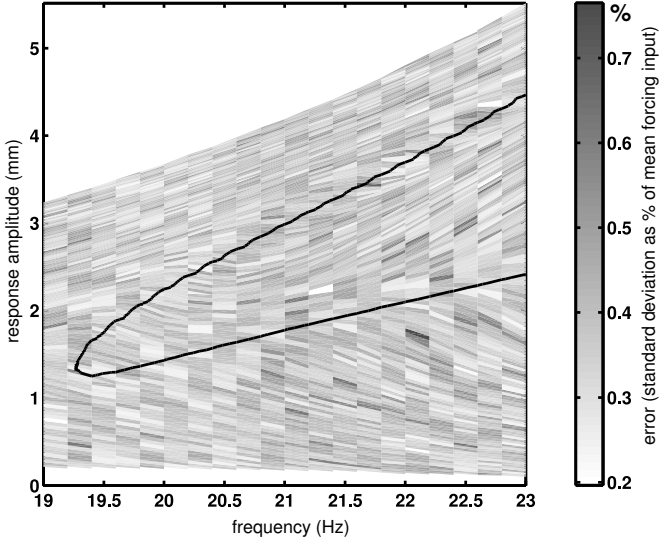


David Barton's results (Bristol)



David Barton's results (Bristol)

Error estimate



Computational example for Hopf bifurcation

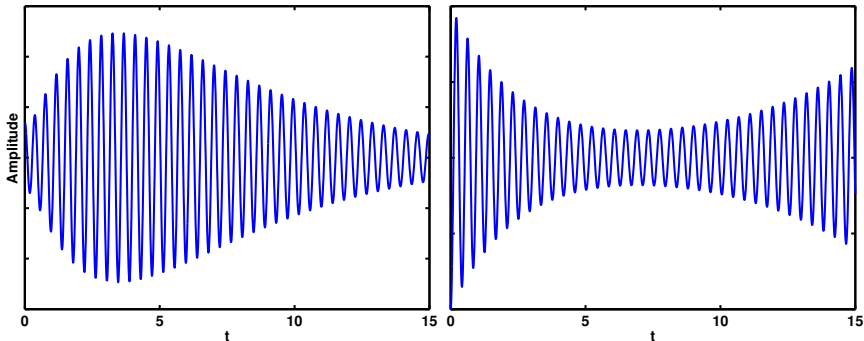
- ▶ Haken-Kelso-Bunz (HKB) oscillator:

$$\ddot{x} + \dot{x}(\alpha x^2 + \beta \dot{x}^2 - \gamma) + \omega^2 x = u + \text{"noise"}$$

- ▶ nonlinearity:
mixture of nonlinear damping terms: $\dot{x}(\alpha x^2 + \beta \dot{x}^2)$
- ▶ Equilibrium at $x = 0$, bifurcation parameter γ
 - ▶ stable for $\gamma < 0$,
 - ▶ unstable for $\gamma > 0$,
 - ▶ Hopf bifurcation at $\gamma = 0$
- ▶ fix $\alpha = 1$, $\omega = 2$,
- ▶ $\beta = -0.1$ (small periodic orbits unstable), or
 $\beta = 0.2$ (small periodic orbits stable)

General problems for parameter scan

- ▶ Do transient oscillations decay or grow (slowly)?
- ▶ Non-normality: below are two linear systems



Effect of noise/disturbance in parameter scan

- ▶ stable Hopf bifurcation:
noise-induced fluctuations gradually grow in amplitude and become more coherent
- ▶ unstable Hopf bifurcation:
same, but at some point before Hopf bifurcation escape occurs.

Matlab demo

Hopf bifurcation as regular root

- ▶ assume that we have entire state as output (x, \dot{x})
- ▶ set (PD control, stabilising for k_1, k_2 large)

$$u = -k_1(x - x_{\text{ref}}(t)) - k_2(\dot{x} - \dot{x}_{\text{ref}}(t)) \quad \text{where}$$

$$x_{\text{ref}}(t) = \delta \sin(2\pi t/T) \quad (\delta \text{ small}).$$

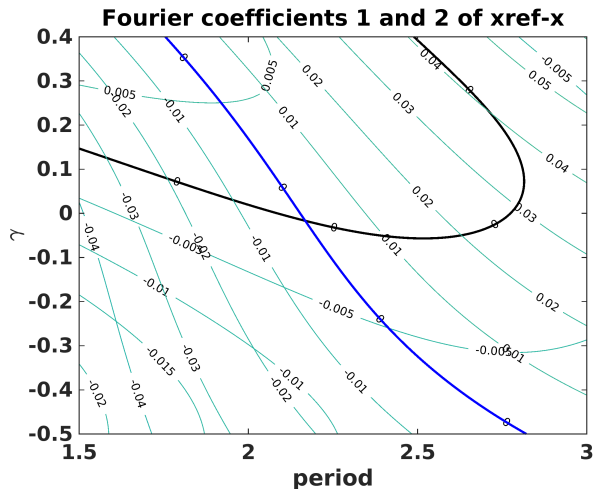
- ▶ Check for $x(t)$ after transients have settled:

$$f_1(T, \gamma) = \int_0^T \sin(2\pi t/T)[x(t) - x_{\text{ref}}(t)] dt$$

$$f_2(T, \gamma) = \int_0^T \cos(2\pi t/T)[x(t) - x_{\text{ref}}(t)] dt$$

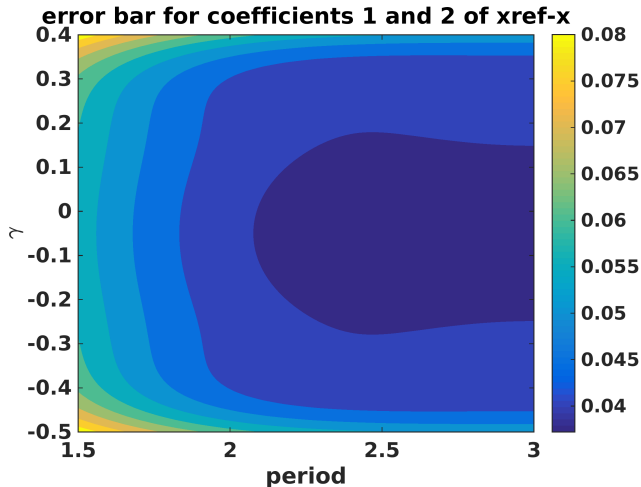
- ▶ If $f_{1,2}(T, \gamma)$ has regular root \Rightarrow uncontrolled system has periodic orbit of amplitude δ , period T at parameter γ .

Hopf bifurcation as regular root



⇒ small-amplitude periodic orbit present

Hopf bifurcation as regular root



⇒ small-amplitude periodic orbit present

Hopf bifurcation as regular root

matlab picture

Comments:

- ▶ transients always finite
- ▶ result depends on gains
becomes independent of gains as noise level $\rightarrow 0$.
- ▶ Hopf bifurcation defined in limit noise level $\rightarrow 0$.

Further remarks about continuation in experiments

Compared to numerical continuation

1. evaluation of $F : (x_{\text{ref}}, p) \rightarrow x$ is slow

⇒ restricted to low dimension of control inputs and small number of (eg) Fourier modes

2. low accuracy of F (relative error $\approx 10^{-2}$ in very clean experiments)

⇒ restricted to well-conditioned problems

⇒ F. Schilder's `coco` toolbox `continex`

3. limiting factor: ability to provide stabilising (!) real-time feedback

⇒ hardest part is problem specific

⇒ new algorithms needed