ON ALGEBRAIC TORSION FORMS AND THEIR SPIN HOLONOMY ALGEBRAS

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ABSTRACT. We study holonomy algebras generated by an algebraic element of the Clifford algebra, or equivalently, the holonomy algebras of certain spin connections in flat space. We provide some series of examples in arbitrary dimensions and prove some general properties of the holonomy algebras under some mild conditions on the generating element. We show that the first non-standard situation to look at appears in dimension 8 and concerns 4-forms. In this case complete structure results are obtained when moreover assuming the 4-form to be self-dual.

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1. Introduction

Let (M^n, g) be a Riemannian spin manifold, with spinor bundle to be denoted by \mathcal{S} . For any differential form T on M, not necessarily of pure degree, one can form

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the linear connection ∇^T on \mathcal{S} by setting

$$\nabla_X^T \psi = \nabla_X \psi + (X \perp T) \psi$$

whenever ψ belongs to $\Gamma(\mathcal{S})$ and X is in TM. Here ∇ is the connection induced by the Levi-Civita connection on the spinor bundle \mathcal{S} . This can be thought of as the spin analogue of a connection with non-trivial torsion on the tangent bundle of M. A special case is when T is actually a 3-form making that ∇^T is the spin connection of a genuine connection with torsion on TM. In low significant dimensions, ranging from 6 to 8 parallel spinors w.r.t to a connection with 3-form torsion are nowadays rather well understood in terms of geometric structures on the tangent bundle to the manifold [4, 10] and extensive effort toward their classification has been made [7, 11, 9]. This has been also studied in connection with the so-called Stromingers type II string equations [15], [3]. Another special case, which no longer reflects the presence of a particular connection at the level of the tangent of M is when T consists of forms of degree 3 and 4, the latter being termed fluxes in physics literature (see [8] and references therein). In all the above mentioned cases one of the issues to understand is under which conditions ∇^T admits parallel spinors, therefore one looks, more generally, at the holonomy representation of ∇^T .

In this paper we shall study the holonomy of the connection ∇^T in the flat case, when moreover T is assumed to have constant coefficients. This is the simplest geometric case one could think of but already raises some interesting and quite difficult algebraic questions. We set

Definition 1.1. Let $(V^n, \langle \cdot, \cdot \rangle)$ be an Euclidean vector space and let $Cl_n(V)$ be its Clifford algebra. Then:

- (i) the fix algebra of T in $Cl_n(V)$ is the Lie-sub-algebra of $Cl_n(V)$ generated by $\{X \perp T : X \in V\}$.
- (ii) the holonomy algebra of some T in $Cl_n(V)$ is given as $\mathfrak{h}_T^* = [\mathfrak{g}_T^*, \mathfrak{g}_T^*]$.

This is motivated by the observation [2] that in the flat case the holonomy algebra of the spin connection ∇^T equals \mathfrak{h}_T^* . When \mathfrak{g}_T^* is *perfect*, that is $\mathfrak{g}_T^* = [\mathfrak{g}_T^*, \mathfrak{g}_T^*]$ the two algebras above coincide and in this respect the fix algebra \mathfrak{g}_T^* appears to be a very useful intermediary object for establishing structure results, although it seems to lack of further geometric content.

Our paper is organised as follows. In section 2 we review a number of elementary facts concerning Clifford algebras and their representations, with accent put on the different phenomenon appearing in some arithmetic series of dimensions. In section 3 we start our study of holonomy algebras by determining - under some mild assumptions on the generating element - the model algebra those are contained in. This allows proving a number of general properties like semisimplicity which appear to be useful when looking at the center of the holonomy algebras. Further on, we investigate the space of the so-called fixed spinors which for some T in $Cl_n(V)$ is defined as

$$Z_T = \{ \psi \in \mathcal{S} : (X \sqcup T)\psi = 0 \text{ for all } X \text{ in } V \}$$

where \mathcal{S} is an irreducible $Cl_n(V)$ module. We provide first order information about these spaces and also discuss some simple examples. The section ends with giving a necessary condition for certain holonomy algebras to be perfect, namely

Theorem 1.1. Let T in $Cl_n^0 \cap Cl_n^+$, where $n \equiv 0 \pmod{4}$ satisfy $T^t = T$. If $Z_T = (0)$ then \mathfrak{g}_T^* is a perfect Lie algebra.

Section 4 describes situations where the holonomy algebras can be directly computed and provides series of useful, in hindsight, examples. Elements of the Clifford algebra which are being looked at are unipotent and squares of spinors, which actually give idempotents. In the latter situation, the dimension (mod 8) of the underlying vector space appears to lead to very different results. More precisely

Theorem 1.2. Let V be an Euclidean vector space with volume form ν and let T belong to $Cl_n(V)$. Then:

- (i) if $n \equiv 0 \pmod{4}$ and T in $Cl_n^0 \cap Cl_n^+$ is unipotent, that is $T^2 = 1 + \nu$, then its holonomy algebra is isomorphic with $\mathfrak{so}(n,1)$.
- (ii) if T is the square of a spinor when $n \equiv 7 \pmod{8}$, respectively the square of a positive spinor when $n \equiv 0 \pmod{8}$ then the holonomy algebra of T is abelian or isomorphic with $\mathfrak{so}(n,1)$ respectively.

This essentially exploits specific features of the powerful squaring construction for spinors [13], [6]. Section 5 forms the core of the present paper and gives a complete classification of holonomy algebras generated by self-dual 4-forms in dimension 8. More precisely, we show

Theorem 1.3. Let V be any oriented Euclidean vector space of dimension 8 and let T be a self-dual four form on V. The holonomy algebra of T is isomorphic to the Lie algebra $\mathfrak{so}(8, 8 - \dim_{\mathbb{R}} Z_T)$ exception made of the cases when

- (i) T is an unipotent element
- (ii) $dim_{\mathbb{R}}Z_T=6$

where the holonomy algebras are $\mathfrak{so}(8,1)$ and $\mathfrak{so}(6,2)$ respectively.

The theorem is proved, exception made of the case when $dim_{\mathbb{R}}Z_T=6$, by making use, as main tool, of the splitting of the space of two forms which can be obtained once a self-dual four form is given combined with the observation that raising the generating element to any odd power leaves the initial holonomy algebra unchanged. In the last section of the paper, we treat by direct computation the special case appearing in (ii) of Theorem 1.3 using the one to one correspondence [12] between the existence of a such a form and that of an SU(4)-structure on our vector space. The paper ends by an appendix, containing the elementary though lengthy proof of a technical Lemma.

2. Preliminaries

In this section we shall recall a number of facts concerning Clifford algebras and spinors, we shall constantly make use of in what follows. A thorough account on all theses notions can be found in [13].

2.1. Clifford algebras. Let V be an n-dimensional vector space over \mathbb{R} equipped with a scalar product, to be denoted by $\langle \cdot, \cdot \rangle$. We shall denote by $Cl_n(V)$ the Clifford algebra associated with $(V, \langle \cdot, \cdot \rangle)$, and if there is no ambiguity on the vector space used we shall simply write Cl_n for $Cl_n(V)$. We recall that $Cl_n(V)$ can be given the

structure of an algebra, with multiplication denoted by " \cdot ": $Cl_n(V) \to Cl_n(V)$ and satisfies

$$(2.1) e \cdot \varphi = e \wedge \varphi - e \, \lrcorner \, \varphi, \varphi \cdot e = (-1)^k (e \wedge \varphi + e \, \lrcorner \, \varphi)$$

whenever e belongs to $\Lambda^1(V)$ and φ is in $\Lambda^k(V)$, although this notation will no longer be used in what follows. Here and henceforth we will identify 1-forms and vectors via the given scalar product $\langle \cdot, \cdot \rangle$. There is a canonical isomorphism between the space $\Lambda^*(V)$ of forms in V and the Clifford algebra $Cl_n(V)$ having the following property. Let $L: Cl_n(V) \to Cl_n(V)$ be defined by

$$L(\varphi) = \sum_{i=1}^{n} e_i \, \varphi \, e_i,$$

whenever $\varphi \in Cl_n(V)$ and for some orthonormal basis $\{e_i\}$, $1 \le i \le p$ in V. Then the eigenspaces of L are the canonical images of $\Lambda^k(V)$:

(2.2)
$$L|_{\Lambda^k} = (-1)^k (2k - n) 1_{\Lambda^k}.$$

Any Clifford algebra comes with two involutions, the first being the transposition map $()^t: Cl_n(V) \to Cl_n(V)$ defined by

$$(e_1 e_2 \dots e_{k-1} e_k)^t = e_k e_{k-1} \dots e_2 e_1,$$

for some orthonormal frame $\{e_i, 1 \leq i \leq n\}$. Note however that the transpose is frame independent and therefore extends to an anti-automorphism of $Cl_n(V)$, i.e.

$$(2.3) \qquad (\varphi_1 \, \varphi_2)^t = \varphi_2^t \, \varphi_1^t$$

for all φ_1, φ_2 in $Cl_n(V)$. The second involution $\alpha : Cl_n(V) \to Cl_n(V)$ results from extending -1_V to an automorphism of the algebra $Cl_n(V)$, in the sense that

$$\alpha(\varphi_1 \, \varphi_2) = \alpha(\varphi_1) \, \alpha(\varphi_2),$$

where φ_1, φ_2 are in $Cl_n(V)$. Since α is an involution it can also be used to obtain a splitting

$$Cl_n(V) = Cl_n^0(V) \oplus Cl_n^1(V)$$

into the \pm -eigenspaces of α . Note this corresponds to the splitting of $\Lambda^*(V)$ into even and respectively odd degree forms. The vector space $Cl_n(V)$ inherits from V a scalar product, still to be denoted by $\langle \cdot, \cdot \rangle$ and having the property that

(2.4)
$$\langle q\varphi_1, \varphi_2 \rangle = \langle \varphi_1, \alpha(q^t)\varphi_2 \rangle$$

$$\langle \varphi_1 q, \varphi_2 \rangle = \langle \varphi_1, \varphi_2 \alpha(q^t) \rangle$$

whenever q, φ_1, φ_2 belong to $Cl_n(V)$. Let now assume that V is oriented by ν in $\Lambda^n(V)$ such that for an oriented frame $\{e_k, 1 \leq k \leq n\}$ this is given as $\nu = e_1 \ldots e_n$. Then it is easy to check that

(2.5)
$$\nu^2 = (-1)^{\frac{n(n+1)}{2}}, \qquad \nu^t = (-1)^{\frac{n(n-1)}{2}}\nu.$$

Now the Hodge star operator $*: \Lambda^k(V) \to \Lambda^{n-k}(V)$ is defined by $\alpha \wedge *\beta = \langle \alpha, \beta \rangle \nu$, for all α, β in $\Lambda^*(V)$ and relates to Clifford multiplication with ν by

(2.6)
$$*\varphi = (-1)^{\frac{k}{2}(k+1)}\varphi\nu = (-1)^{\frac{k}{2}(2n-k+1)}\nu\varphi,$$

for φ in $\Lambda^k(V) \subset Cl_n(V)$. Moreover, we have

(2.7)
$$\varphi \nu = \nu \varphi, \quad \text{for all } \varphi \text{ in } Cl_n^0,$$

$$\varphi \nu = (-1)^{n+1} \nu \varphi, \quad \text{for all } \varphi \text{ in } Cl_n^1.$$

In particular, when $n \equiv 1 \pmod{2}$, the volume element ν belongs to the center of the Clifford algebra Cl_n .

If $n \equiv 0 \pmod{4}$ then $\nu^2 = 1$ whence the Hodge star operator, realised as in (2.6), provides a decomposition of the Clifford algebra into self-dual and anti-self-dual elements:

$$(2.8) Cl_n(V) = Cl_n^+(V) \oplus Cl_n^-(V),$$

where $\nu\varphi = \pm \varphi$ whenever φ belongs to $Cl_n^{\pm}(V)$.

2.2. The space of spinors. We need to recall some elementary facts about spinors. Let \mathcal{S} be an irreducible $Cl_n(V)$ -module. We shall call \mathcal{S} the space of spinors and elements $\psi \in \mathcal{S}$ spinors and denote by $\mu : Cl_n(V) \to \operatorname{End}(\mathcal{S})$ the Clifford multiplication acting on \mathcal{S} . On \mathcal{S} we have the usual scalar product $\langle \psi_1, \psi_2 \rangle$ for two spinors $\psi_1, \psi_2 \in \mathcal{S}$ which has the following property

(2.9)
$$\langle \varphi \psi_1, \psi_2 \rangle = \langle \psi_1, \alpha(\varphi^t) \psi_2 \rangle,$$

for all φ in Cl_n and ψ_1, ψ_2 in \mathcal{S} . Recalling that

(2.10)
$$\alpha(\varphi^t) = (-1)^{\frac{k}{2}(k+1)} \varphi, \qquad \varphi \in \Lambda^k(V) \subset Cl_n(V).$$

it follows that the Clifford multiplication operator $\mu_{\varphi}: \mathcal{F} \to \mathcal{F}$, with an element φ of $\Lambda^k(V)$, is symmetric when $k \equiv 0, 3 \pmod{4}$ and anti-symmetric when $k \equiv 1, 2 \pmod{4}$. Now when $n \equiv 0, 3 \pmod{4}$ the volume form ν squares to 1. If $n \equiv 0 \pmod{4}$ this allows splitting the irreducible, real Clifford \mathcal{F} as $\mathcal{F} = \mathcal{F}^+ \oplus \mathcal{F}^-$, where ν acts as ± 1 on \mathcal{F}^\pm . Peculiar to the case when $n \equiv 3 \pmod{4}$ is the fact that any irreducible, real, Clifford module \mathcal{F} has either $\nu\psi = -\psi$ for all ψ in \mathcal{F} or $\nu\psi = \psi$ for all ψ in \mathcal{F} . Both possibilities can occur and produce different Cl_n -representations. As a convention, in what follows we shall always work with the latter representation. Let us also mention that when $n \equiv 3 \pmod{4}$ we have that $\nu\varphi = \varphi\nu$ for all φ in Cl_n and that α interchanges Cl_n^+ and Cl_n^- , that is realises an isomorphism $\alpha: Cl_n^+ \to Cl_n^-$. For it will be used constantly in what follows we also recall the following stability Lemma.

Lemma 2.1. The following stability conditions hold:

(i) when $n \equiv 0 \pmod{4}$,

$$\begin{split} \varphi \mathcal{S}^+ &\subseteq \mathcal{S}^+, \quad \forall \varphi \in Cl_n^0(V) \cap Cl_n^+(V) \\ \varphi \mathcal{S}^+ &= 0, \quad \forall \varphi \in Cl_n^0(V) \cap Cl_n^-(V) \\ \varphi \mathcal{S}^+ &= 0, \quad \forall \varphi \in Cl_n^1(V) \cap Cl_n^+(V) \\ \varphi \mathcal{S}^+ &\subseteq \mathcal{S}^+, \quad \forall \varphi \in Cl_n^1(V) \cap Cl_n^-(V), \end{split}$$

(ii) while for $n \equiv 3 \pmod{4}$ we have $Cl_n^+ \mathcal{S} \subseteq \mathcal{S}, Cl_n^- \mathcal{S} = 0$.

The proof, which is left to the reader, follows from the above properties of Clifford multiplication with ν . Similar statements can be easily made on \mathcal{F}^- when $n \equiv 0 \pmod{4}$. We end this section by recalling two more well known facts, with proofs given for the sake of completeness.

Lemma 2.2. If ζ in Cl_n satisfies $[\zeta, \Lambda^2(V)] = 0$ then ζ belongs to $(1, \nu)$.

Proof. Since $\zeta XY = XY\zeta$ for all X,Y in V it follows that $Y(X\zeta X)Y = |X|^2|Y|^2\zeta$ and further, after a double tracing $L^2\zeta = n^2\zeta$. The eigenvalues of L^2 being $(2p-n)^2$ it follows that the only degrees present in ζ are 0 and n. Therefore ζ is a linear combination of 1 and ν .

Lemma 2.3. Let \mathcal{F} be any irreducible, real, Cl_n module where $n \equiv 0, 3 \pmod{4}$. For any φ in Cl_n such that $\alpha(\varphi^t) = \varphi$ we have

- (i) $Tr(\mu_{\varphi}) = dim_{\mathbb{R}} \mathcal{S} \langle \varphi, 1 \rangle$, if $n \equiv 0 \pmod{4}$,
- (ii) $Tr(\mu_{\varphi}) = 2 \dim_{\mathbb{R}} \mathcal{S}(\varphi, 1)$, if $n \equiv 3 \pmod{4}$.

Proof. Let us consider the linear sub-space of Cl_n given by $\mathcal{S} = \{\varphi : \alpha(\varphi^t) = \varphi\}$. We now pick some orthonormal basis $\{e_i : 1 \leq i \leq n\}$ in V and observe that $e^{ij}\mathcal{S}e^{ij}\subseteq\mathcal{S}$, where $e^{ij}=e^i\wedge e^j, i\neq j$. Let now $t:\mathcal{S}\to\mathbb{R}$ be given as $t(\varphi)=Tr(\mu_\varphi)$ for all φ in \mathcal{S} . Since this is linear, it can be written as $t=\langle\cdot,T\rangle$ for some T in \mathcal{S} . From the independence of the trace of some orthonormal basis in \mathcal{S} and (2.9) we get $t(e^{ij}\varphi e^{ij})=-t(\varphi)$ for all φ in \mathcal{S} . Using (2.4) this results in having $e^{ij}Te^{ij}=-T$ or further $[T,e^{ij}]=0$ for all $1\leq i\neq j\leq n$, where we have used that $(e^{ij})^2=-1$. Henceforth $[T,\Lambda^2]=0$, leading by Lemma 2.2 to $T=\lambda_1+\lambda_2\nu$ for some λ_1,λ_2 in \mathbb{R} which can be computed as $\lambda_1=Tr(\mu_1)=dim_{\mathbb{R}}\mathcal{S}$ and $\lambda_2=Tr(\mu_\nu)=dim_{\mathbb{R}}\mathcal{S}^+-dim_{\mathbb{R}}\mathcal{S}^-=0$ for $n\equiv 0\pmod 4$ and accordingly $Tr(\mu_\nu)=dim_{\mathbb{R}}\mathcal{S}$ for $n\equiv 3\pmod 4$, where we have used that ν belongs to \mathcal{S} .

3. Structure results

3.1. The general setup. In this section our aim is mainly to locate some classes of holonomy algebras inside the Clifford algebra and derive a number of general properties they must satisfy. Let A be the subset of Cl_n given by

$$A = \{ \varphi \in Cl_n : \varphi^t = -\varphi \}.$$

This is meant to be the model algebra for most classes of holonomy algebras we will be looking at, in a sense to be made precise below.

Lemma 3.1. The following hold:

- (i) A is a Lie sub-algebra of $(Cl_n, [\cdot, \cdot])$.
- (ii) The symmetric bilinear form $\beta(\varphi_1, \varphi_2) = \langle \varphi_1, \alpha(\varphi_2) \rangle$ is non-degenerate on A and invariant, that is

$$\beta([\varphi_1, \varphi_2], \varphi_3) = -\beta([\varphi_1, \varphi_3], \varphi_2)$$

whenever $\varphi_k, 1 \leq k \leq 3$ belong to A.

Proof. (i) Follows immediately from anti-symmetrising that $(\varphi_1\varphi_2)^t = \varphi_2^t\varphi_1^t = \varphi_2\varphi_1$ whenever φ_1, φ_2 belong to A.

(ii) The non-degeneracy of β follows from A being preserved by the involution α . Now

$$\beta([\varphi_1, \varphi_2], \varphi_3) = \langle [\varphi_1, \varphi_2], \alpha(\varphi_3) \rangle = \langle \varphi_1 \varphi_2 - \varphi_2 \varphi_1, \alpha(\varphi_3) \rangle$$
$$= \langle \varphi_2, \alpha(\varphi_1^t) \alpha(\varphi_3) - \alpha(\varphi_3) \alpha(\varphi_1^t) \rangle$$
$$= \langle \varphi_2, \alpha([\varphi_1, \varphi_3]) \rangle = -\beta(\varphi_2, [\varphi_1, \varphi_3]).$$

Since A is stable under α , it inherits from Cl_n a bi-grading $A = A^0 \oplus A^1$ into its even respectively its odd degree components. The usual rules $[A^0, A^0] \subseteq A^0$, $[A^0, A^1] \subseteq A^1$, $[A^1, A^1] \subseteq A^0$ apply, in particular A^0 is a Lie sub-algebra of A. We can obtain now first order information about some of the holonomy algebras, by assuming the generating element to be well related to the standard decompositions of Cl_n .

Proposition 3.1. Let T belong to Cl_n^0 and satisfy $T^t = T$. Then:

- (i) \mathfrak{g}_T^* is a Lie sub-algebra of A
- (ii) $\alpha(\mathfrak{g}_T^*) = \mathfrak{g}_T^*$.

Proof. (i) follows eventually after checking that the generating set $\{X \sqcup T : X \in V\}$ is contained in A as

$$(X \, \lrcorner \, T)^t = -X \, \lrcorner \, T^t = -X \, \lrcorner \, T$$

for all X in V. To prove (ii) we notice that $\alpha(X \, \lrcorner \, T) = -X \, \lrcorner \, T$ for all X in V, in other words α preserves the generating set. Since α is a Lie algebra automorphism preserving A, the result follows.

Therefore, for any T in Cl_n , the Lie algebra \mathfrak{g}_T^* splits as

$$\mathfrak{g}_T^*=\mathfrak{g}_T^{*,0}\oplus\mathfrak{g}_T^{*,1},$$

where the obvious notation applies. Getting closer to the specific features of the algebras \mathfrak{g}_T^* requires some Lie algebra background we shall now briefly outline. For our setup most convenient is to adopt the following

Definition 3.1. Let \mathfrak{g} be a real Lie algebra. It is called semisimple if it admits a symmetric bilinear form β which is non-degenerate and satisfies

$$\beta([\varphi_1, \varphi_2], \varphi_3) = -\beta([\varphi_1, \varphi_3], \varphi_2)$$

for all $\varphi_1, \varphi_2, \varphi_3$ in \mathfrak{g} .

This essentially ensures that any ideal \mathfrak{i} of \mathfrak{g} has trivial extension, that is there exists \mathfrak{i}^{\perp} such that $\mathfrak{g} = \mathfrak{i} \oplus \mathfrak{i}^{\perp}$, where \mathfrak{i}^{\perp} denotes the orthogonal complement of \mathfrak{i} w.r.t the non-degenerate form β . In particular

Proposition 3.2. Let \mathfrak{g} be a real Lie algebra. If \mathfrak{g} is semisimple then $[\mathfrak{g},\mathfrak{g}] = \mathfrak{g}$ if and only if it has trivial center. Here the center $Z(\mathfrak{g})$ of \mathfrak{g} is given as $Z(\mathfrak{g}) = \{\zeta \in \mathfrak{g} : [\zeta,\mathfrak{g}] = 0\}$.

Real Lie algebras \mathfrak{g} satisfying $[\mathfrak{g},\mathfrak{g}] = \mathfrak{g}$ are termed *perfect* and Proposition 3.2 provides a criterion for checking this, to be used later on.

Lemma 3.2. Suppose that $n \equiv 0 \pmod{8}$. The following hold:

- (i) A is a semi-simple Lie algebra.
- (ii) A isomorphic to $\mathfrak{so}(d,d)$ where $d=\frac{1}{2}dim_{\mathbb{R}}\mathcal{S}$ and \mathcal{S} is the irreducible real Cl_n module.
- (iii) the adjoint representation of A^0 on A^1 is irreducible.

Proof. (i) follows from the non-degeneracy of β on A, which is due to $\alpha(A) = A$. (ii) Let us equip \mathcal{S} with the scalar product $\hat{\beta}$ which leaves \mathcal{S}^+ and \mathcal{S}^- orthogonal and equals $\pm \langle \cdot, \cdot \rangle$ on \mathcal{S}^{\pm} . In short, $\hat{\beta}(x,y) = \langle \nu x, y \rangle$ for all x, y in \mathcal{S} . If φ is in A it is easy to check that $\hat{\beta}(\mu_{\varphi}x,y) + \hat{\beta}(\mu_{\varphi}y,x) = 0$ for all x,y in \mathcal{S} , that is μ_{φ} belongs to $\mathfrak{so}(\mathcal{S},\hat{\beta}) \cong \mathfrak{so}(d,d)$. But when $n \equiv 0 \pmod{8}$ the Clifford multiplication gives a linear isomorphism $\mu: Cl_n \to End(\mathcal{S},\mathcal{S})$ which is also a Lie algebra isomorphism

(iii) follows standardly from (ii).
$$\Box$$

Similar results can be proved in the remaining series of dimensions but this is somewhat beyond the scope of the present paper. In the same vein

Proposition 3.3. For any T in Cl_n^0 with $T^t = T$ the Lie algebra \mathfrak{g}_T^* is semisimple.

Proof. We need only see that the restriction of β to \mathfrak{g}_T^* is non-degenerate. But this follows easily from the fact that α preserves \mathfrak{g}_T^* .

Lemma 3.3. Let T be in Cl_n^0 . Then

$$\mathfrak{g}_{eTe}^* = e\,\mathfrak{g}_T^*\,e$$

for any unit vector e in V.

and our claim follows.

Proof. At first we notice that eTe still belongs to Cl_n^0 . We have

$$-2X \mathrel{\lrcorner} (eTe) = X(eTe) - (eTe)X = -e(XT)e + e(TX)e = 2e(X \mathrel{\lrcorner} T)e$$

for all X in $(e)^{\perp}$. Similarly, $e \perp (eTe) = e(e \perp T)e$ hence $X \perp (eTe) = e(F_eX \perp T)e$ for all X in V, where F_e is the invertible endomorphism of V which equals -1 on $(e)^{\perp}$ and 1 on (e). Let $\rho_e: Cl_n \to Cl_n$ be defined as $\rho_e(\varphi) = e\varphi e$ for all φ in Cl_n^0 . It therefore maps the generating set of \mathfrak{g}_T^* onto that of $e\mathfrak{g}_T^*e$ and since $-\rho_e$ is a Lie algebra isomorphism it is easy to conclude.

When $n \equiv 0 \pmod{4}$ the map ρ_e intertwines, say, $Cl_n^+ \cap Cl_n^0$ and $Cl_n^- \cap Cl_n^0$, therefore from the Lemma above we see that the holonomy algebra does not distinguish between the generating elements being in Cl_n^+ or Cl_n^- . Hence all results obtained for holonomy algebras generated by elements in Cl_n^+ extend automatically to generating elements in Cl_n^- . We end this section by an example of forms when the holonomy algebras can be easily computed.

Proposition 3.4. Let $(V^n, \langle \cdot, \cdot \rangle)$ be an Euclidean vector space oriented by ν in $\Lambda^n(V)$. Then:

- (i) $\mathfrak{g}_{\nu}^* = \mathfrak{so}(n,1)$ for $n \equiv 0,1 \pmod{4}$
- (ii) $\mathfrak{g}_{\nu}^* = \mathfrak{so}(n+1)$ for $n \equiv 2, 3 \pmod{4}$ (iii) in both cases $\mathfrak{g}_{\nu}^* \psi = 0$ if and only if $\psi = 0$, for any $\psi \in \mathcal{S}$.

Proof. Let us first notice that the generating set $\{X \sqcup \nu : X \in V\}$ is isomorphic to V, since the volume form ν is non-degenerate. Keeping in mind that by (2.1) we have $X \perp \nu = -X\nu$ and using (2.5), (2.7) this yields

$$[X \, \lrcorner \, \nu, Y \, \lrcorner \, \nu] = (-1)^{\frac{1}{2}(n+1)(n+2)} [X, Y] = 2(-1)^{\frac{1}{2}(n+1)(n+2)} X \wedge Y$$

for all X, Y in V. Similarly we get for the triple commutators

$$[\alpha, X \rfloor \nu] = [X, \alpha]\nu = 2 FX \rfloor \nu$$

for all X in V, where $\alpha = \langle F \cdot, \cdot \rangle$ belongs to $\Lambda^2(V)$. Therefore $\mathfrak{g}_{\nu}^* = V \oplus \Lambda^2(V)$ as a vector space and the claims in (i) and (ii) follow from the commutator rules above. (iii) follows easily from the invertibility of ν in Cl_n , as defined in Def. 3.2.

3.2. The set of fixed spinors. As it will appear below the holonomy algebra of some element T in Cl_n is intimately related to the space of spinors fixed by T, which we recall to be defined as

$$(3.1) Z_T = \{ \psi \in \mathcal{S} : (X \perp T)\psi = 0, \text{ for all } X \in V \}.$$

Notice that if (the non-zero) T is of degree 1 or 2 the set Z_T is obviously reduced to zero and moreover the latter holds for forms of degree 3 (see [2]). We now gather a number of basic facts concerning the set Z_T . If $n \equiv 0 \pmod{4}$ we split Z_T along the splitting $\mathcal{S} = \mathcal{S}^+ \oplus \mathcal{S}^-$ and get

$$Z_T = Z_T^+ \oplus Z_T^-$$

where the obvious notation applies.

Lemma 3.4. Let T belong to $Cl_n^0 \cap Cl_n^+$ where $n \equiv 0 \pmod{4}$. Then

(i)

$$Z_T^+ = \{ \psi \in \mathcal{S}^+ : T\psi = 0 \}$$

$$Z_T^- = \{ \psi \in \mathcal{S}^- : TV\psi = 0 \}$$

(ii) if n = 8, then $Z_T^- = (0)$ provided T does not vanish.

Proof. (i) follows directly from the stability conditions.

(ii) If $Z_T^- \neq (0)$ there exists a non-zero ψ in \mathcal{F}^- with $TV\psi = 0$. But in 8-dimensions $V\psi = \mathcal{F}^+$ hence $T\mathcal{F}^+ = 0$ whence T = 0, a contradiction.

Definition 3.2. An element T of $Cl_n^0 \cap Cl_n^+$ is invertible in $Cl_n/(1,\nu)$ if there exists T^{-1} in $Cl_n^0 \cap Cl_n^+$ with $T^{-1}T = TT^{-1} = 1 + \nu$.

Actually a necessary and sufficient condition for Z_T to vanish is

Proposition 3.5. Let T belong to $Cl_n^0 \cap Cl_n^+$, $n \equiv 0 \pmod{4}$ such that $T^t = T$. Then $Z_T = (0)$ if and only if T is invertible in $Cl_n/(1, \nu)$.

Proof. If T is invertible Lemma 3.4 yields immediately the vanishing of Z_T^{\pm} hence that of T. Suppose now that $Z_T = (0)$ and let $L_T : Cl_n^+ \cap Cl_n^0 \to Cl_n^+ \cap Cl_n^0$ be left multiplication with T. If φ is in the kernel of L_T it follows that $T(\varphi \mathcal{S}^+) = 0$ and moreover, since $\varphi \mathcal{S}^+ \subseteq \mathcal{S}^+$ Lemma 3.4 tells us that $\varphi \mathcal{S}^+ \subseteq Z_T$ hence $\varphi \mathcal{S}^+ = 0$. Therefore φ vanishes and it follows that L_T is injective, thus invertible and this provides easily the required inverse for T in $Cl_n/(1,\nu)$ given that $1_{Cl_n^+ \cap Cl_n^0} = \frac{1}{2}L_{1+\nu}$.

In the rest of this section we shall present examples of situations when the set of fixed spinor can be seen directly to be trivial.

Proposition 3.6. Let α in $\Lambda^2(V)$ be a two-form such that $T = \alpha \wedge \alpha \neq 0$. Then $Z_T = (0)$.

Proof. Let F be the skew-symmetric endomorphism associated to α via the metric g, that is $\alpha = \langle F \cdot, \cdot \rangle$. We have $X \sqcup (\alpha \wedge \alpha) = 2(X \sqcup \alpha) \wedge \alpha = 2FX \wedge \alpha$ for all X in V. Let now ψ be in Z_T and let us set $r = \operatorname{rank}(F)$. From

$$(X \rfloor T)\psi = 2 (FX \land \alpha)\psi = 0$$

follows $(X \wedge \alpha)\psi = 0$ for all X in Im(F). By Clifford contraction we get immediately

$$\sum_{e_i \in Im(F)} e_i(e_i \wedge \alpha)\psi = -\sum_{e_i \in Im(F)} (e_i \sqcup (e_i \wedge \alpha))\psi = (2-r) \alpha\psi = 0,$$

which leads to $\alpha \psi = 0$ since having r = 2 would imply $T = \alpha \wedge \alpha = 0$, a contradiction. Therefore

$$0 = X\alpha\psi = (X \wedge \alpha - X \, \lrcorner \, \alpha)\psi$$

for all X in V and since $(X \wedge \alpha)\psi = 0$ for all X in Im(F) we are lead to $(X \sqcup \alpha)\psi = 0$ for all X in Im(F). It follows that $\psi = 0$ as $\alpha \neq 0$.

A very simple observation, which appears to be useful in low dimensions is

Lemma 3.5. Let T belong to $\Lambda^k(V)$, $k \neq 0$. Then:

- (i) $TZ_T = 0$
- (ii) $Z_T = Z_{\nu T}$

Proof. (i) If ψ belongs to Z_T we have $(X \perp T)\psi = 0$ for all X in V. Therefore $\sum_{i=1}^{n} e_i(e_i \perp T)\psi = 0$ for some orthonormal frame $\{e_i, 1 \leq i \leq n\}$ leading to $kT\psi = 0$ and the claim follows.

(ii) From (i) we get that $Z_T = \{ \psi \in \mathcal{S} : TV\psi = 0 \}$ for any pure degree form T and the claim follows easily.

Proposition 3.7. For any T in $\Lambda^k(\mathbb{R}^n)$, $n \leq 7$ the set Z_T is trivial.

Proof. This is obvious when k=1,2 and when k=3 it was proved in [2]. Now if $k \geq 4$ we have $Z_T = Z_{\nu T}$ and since νT has degree $n-k \leq 3$ we conclude by the above.

Therefore the first case of interest is that of dimension 8, which will be studied in detail in the latter part of the paper.

3.3. **Perfect holonomy algebras.** In this section we shall examine situations when the holonomy algebra \mathfrak{g}_T^* for some T in Cl_n is perfect. We will see this is not always the case since those can be abelian by the examples in the next section. However we will show that it is possible to give necessary conditions to that extent. Let us set first a preparatory Lemma.

Lemma 3.6. If ζ_1, ζ_2 in Cl_n satisfy $\zeta_1 X = X\zeta_2$ for all X in V then ζ_1 and ζ_2 belong to $(1, \nu)$.

Proof. It follows that $-|X|^2\zeta_1 = X\zeta_2X$ for all X in V and further

$$|X|^2 Y \zeta_2 Y = |Y|^2 X \zeta_2 X$$

for all X, Y in V. By left multiplication with some non-zero X we get $(XY)\zeta_2Y = -|Y|^2\zeta_2X$ and now right multiplication with a non-zero Y yields $(XY)\zeta_2 = \zeta_2(XY)$

for all X, Y in V. Now ζ_2 is in $(1, \nu)$ by Lemma 2.2 and it is easy to see this implies the claim for ζ_1 as well.

Proposition 3.8. Suppose that $n \equiv 0 \pmod{4}$ and let T satisfy $\nu T = T\nu = T$ and $T^t = T$. If $Z_T = (0)$ then \mathfrak{g}_T^* has trivial center.

Proof. If ζ in $Z(\mathfrak{g}_T^*)$ we must clearly have

$$[\zeta, XT - TX] = 0$$

for all X in V. Since XT-TX belongs to Cl_n^1 for all X in V by applying α to the equation above we get that $[\alpha(\zeta), XT-TX] = 0$ for all X in V, hence after splitting ζ into its even resp. odd components it is enough to treat (3.2) when ζ belongs to Cl_n^0 resp. Cl_n^1 .

Case I: ζ belongs to Cl_n^0

Since $\nu(XT-TX) = -(XT+TX)$ for all X in V and $\nu\zeta = \zeta\nu$ after left multiplication of (3.2) with the volume form we get $\zeta(XT+TX) = (XT+TX)\zeta$ whenever X belongs to V. Taking linear combinations with (3.2) gives further

$$\zeta XT = XT\zeta$$
$$\zeta TX = TX\zeta$$

for all X in V. Since $Z_T = (0)$ we know that T must be invertible in $Cl_n/(1, \nu)$ (see Proposition 3.5) hence using the second equation above we have

$$(T^{-1}\zeta T)X = T^{-1}TX\zeta = (1+\nu)X\zeta = X(1-\nu)\zeta$$

for all X in V. But from Lemma 3.6, (ii) we get that $(1 - \nu)\zeta$ belongs to $(1, \nu)$ and hence vanishes as $((1 - \nu)\zeta)^t = -(1 - \nu)\zeta$. From the vanishing of $(1 - \nu)\zeta$ it follows that $T^{-1}\zeta T = 0$ and this leads after right multiplication with T^{-1} resp. left multiplication with T to $(1 + \nu)\zeta = 0$. Thus $\zeta = 0$ in this case.

Case II: ζ belongs to Cl_n^1 .

We have as before $\zeta(XT-TX)=(XT-TX)\zeta$ for all X in V. But $\nu(XT-TX)=-(XT+TX)$ for all X in V because $\nu T=T$ and since $\nu \zeta=-\zeta \nu$ after left multiplication with the volume form we obtain

$$\zeta(XT + TX) = -(XT + TX)\zeta$$

for all X in V. Taking linear combinations with the original equation gives now

$$(\zeta T)X = -X(T\zeta)$$
$$\zeta XT = -TX\zeta$$

for all X in V. Using Lemma 3.6, (ii) we then get that ζT and $T\zeta$ belong to $(1,\nu)$ and therefore must vanish since elements of Cl_n^1 . The invertibility of T in $Cl_n/(1,\nu)$ leads then to $(1+\nu)\zeta=0$ whence left multiplication by $1+\nu$ in the second equation above gives $TX\zeta=0$ for all X in V. Again the invertibility of T implies that $(1+\nu)X\zeta=X(1-\nu)\zeta=0$ for all X in V and we conclude that $(1-\nu)\zeta=0$ hence $\zeta=0$ and the proof is finished.

Summarising, after making use of the semisimplicity of \mathfrak{g}_T^* and of Proposition 3.2, we obtain

Theorem 3.1. Let T belong to $Cl_n^0 \cap Cl_n^+$ where $n \equiv 0 \pmod{4}$ and satisfy $T^t = T$. If moreover $Z_T = (0)$, the algebra \mathfrak{g}_T^* is perfect, that is

$$\mathfrak{g}_T^* = [\mathfrak{g}_T^*, \mathfrak{g}_T^*].$$

4. Holonomy algebras from distinguished Clifford algebra elements

4.1. Unipotent elements. In this section we shall compute directly the holonomy algebra of an unipotent element T of Cl_n^+ , $n \equiv 0 \pmod{4}$ as introduced below.

Definition 4.1. Let T belong to Cl_n^+ where $n \equiv 0 \pmod{4}$. It is called unipotent if it satisfies $T^t = T$ and $T^2 = 1 + \nu$.

In particular any unipotent element T belongs to Cl_n^0 . We need first to state and prove the following preliminary result, to be used later on as well.

Lemma 4.1. Let T belong to $Cl_n^+ \cap Cl_n^0$ where $n \equiv 0 \pmod{4}$. Then:

$$4[X \sqcup T, Y \sqcup T] = -T[X, Y]T + YT^{2}X - XT^{2}Y$$

whenever X, Y belong to V.

Proof. Follows directly from the stability relations under the form TXT = 0 for all X in V. Details are left to the reader.

Theorem 4.1. The holonomy algebra of any unipotent element T of Cl_n where $n \equiv 0 \pmod{4}$ is isomorphic to $\mathfrak{so}(n,1)$ and moreover we must have $Z_T = (0)$.

Proof. For notational convenience let $E_T = \{X \rfloor T : X \in V\}$ be the generating set of \mathfrak{g}_T^* . It is isomorphic to V under the map $\iota^1 : V \to E_T, \iota^1(X) = X \rfloor T$. Here only the injectivity of ι^1 has to be proved, and indeed, if $\iota^1(X) = 0$ it follows that XT = TX and further $0 = TXT = T^2X$ leading to the vanishing of X. Now by Lemma 4.1 combined with the unipotency of T the space $[E_T, E_T]$ equals $\{T\alpha T + (1 - \nu)\alpha : \alpha \in \Lambda^2(V)\}$. This is isomorphic to $\Lambda^2(V)$ under $\iota^2 : \Lambda^2(V) \to [E_T, E_T], \iota^2(\alpha) = T\alpha T + (1 - \nu)\alpha$. Indeed, if $\iota^2(\alpha) = 0$ we find that $T\alpha T + (1 - \nu)\alpha = 0$ but then both summands vanish as the first is in Cl_n^+ and the second in Cl_n^- . Hence $\alpha = 0$ and so ι^2 is injective. Moreover

$$[T\alpha T + (1-\nu)\alpha, T\beta T + (1-\nu)\beta] = [T\alpha T, T\beta T] + [(1-\nu)\alpha, (1-\nu)\beta]$$
$$+[(1-\nu)\alpha, T\beta T] + [T\alpha T, (1-\nu)\beta]$$
$$=[T\alpha T, T\beta T] + [(1-\nu)\alpha, (1-\nu)\beta]$$

for all α, β in $\Lambda^2(V)$ after using that T belongs to Cl_n^+ . Obviously

$$[(1 - \nu)\alpha, (1 - \nu)\beta] = (1 - \nu)^2 [\alpha, \beta] = 2(1 - \nu)[\alpha, \beta]$$

and moreover the unipotency of T leads easily to $[T\alpha T, T\beta T] = (1 + \nu)T[\alpha, \beta]T = 2T[\alpha, \beta]T$. Altogether this yields

$$2\iota^2[\alpha,\beta] = [\iota^2\alpha,\iota^2\beta]$$

for all α, β in $\Lambda^2(V)$, in other words $\frac{1}{2}\iota^2: \Lambda^2(V) \to [E_T, E_T]$ is a Lie algebra isomorphism. Now, we compute

$$\begin{aligned} -2[T\alpha T + (1-\nu)\alpha, X \, \lrcorner \, T] = & [T\alpha T + (1-\nu)\alpha, XT - TX] \\ = & [T\alpha T, XT] + [(1-\nu)\alpha, XT] \\ & - [T\alpha T, TX] - [(1-\nu)\alpha, TX] \end{aligned}$$

We now estimate each term separately. We have

$$[T\alpha T, XT] = -XT^2\alpha T = -X(1+\nu)\alpha T = -2X\alpha T$$

after using TXT = 0. Similarly, $[T\alpha T, TX] = -2T\alpha X$ and using furthermore that T belongs to Cl_n^+ we finally obtain

$$-2[T\alpha T + (1-\nu)\alpha, X \rfloor T] = -2X\alpha T + 2\alpha XT - 2T\alpha X + 2TX\alpha.$$

This ends by saying that $[\iota^2\alpha, \iota^1X] = \iota^1[X, \alpha]$ whenever α belongs to $\Lambda^2(V)$ and X in V. Summarising, it follows that $\mathfrak{g}_T^* = E_T \oplus [E_T, E_T]$ and moreover

$$\iota^1 \oplus \frac{1}{2}\iota^2 : V \oplus \Lambda^2(V) \to \mathfrak{g}_T^*$$

realises the desired Lie algebra isomorphism with $\mathfrak{so}(n,1)$. Now since any unipotent element T of Cl_n^+ is clearly invertible in $Cl_n/(1,\nu)$ Proposition 3.5 implies that $Z_T = (0)$.

- **Remark 4.1.** (i) We shall see in the next section that unipotent elements naturally play a special role in the classification of holonomy algebras of self-dual 4-forms in dimension 8.
- (ii) Explicit examples of unipotent elements are easy to make. When the dimension of our vector space V satisfies $\dim_{\mathbb{R}} V \equiv 0 \pmod{4}$ we see that $\frac{1}{\sqrt{2}}(1+\nu)$ is unipotent and therefore Theorem 4.1 recovers partly results in Proposition 3.4. Moreover, if we take V_1, V_2 to be Euclidean vector spaces of dimensions $\equiv 0 \pmod{4}$ oriented by volume forms $\nu_k, k = 1, 2$. Then $\frac{1}{\sqrt{2}}(\nu_1 + \nu_2)$ is an unipotent element of the direct product space $V_1 \times V_2$.
- 4.2. **Squares of spinors.** We shall first recall in what follows some facts about the squaring construction in two series of dimensions. To begin with, let $(V^n, \langle \cdot, \cdot \rangle)$ be an Euclidean vector space which furthermore is supposed to be oriented, with orientation given by ν in $\Lambda^n(V)$. A peculiar property of the Clifford multiplication when $n \equiv 8 \pmod{8}$ is then to give an isomorphism (see [13]):

$$(4.1) \mu: Cl_n \to \operatorname{Hom}_{\mathbb{R}}(\mathcal{S}, \mathcal{S})$$

where \mathcal{S} is the irreducible real Cl_n module. When $n \equiv 7 \pmod{8}$ this still holds provided Cl_n is replaced by Cl_n^+ . Let us now fix a spinor $x \in \mathcal{S}^+$ (or in \mathcal{S} if $n \equiv 7 \pmod{8}$), which we normalise to |x| = 1. Then the isomorphism (4.1) gives rise to an element $x \otimes x \in Cl_n$ (or Cl_n^+ when $n \equiv 7 \pmod{8}$) such that:

$$(4.2) (x \otimes x) \psi = \langle \psi, x \rangle x$$

for all ψ in \mathcal{S} . The element $x \otimes x$ consists of forms of various degrees and is customarily called the square of x. Indeed it is well known [13] that

(4.3)
$$x \otimes x = \sum_{k \equiv 0, 3 \pmod{4}}^{n} (x \otimes x)_{k},$$

where $(x \otimes x)_k$ denotes the projection of $(x \otimes x)$ onto $\Lambda^k(V)$. Note that when $n \equiv 0$ (mod 8) the odd degrees are not present. Below we list some of the properties of $x \otimes x$, of relevance for our study.

Lemma 4.2. Let x be a unit length spinor in \mathcal{S}^+ , where $n \equiv 7,8 \pmod{8}$. The following hold:

- (i) the spinor square $x \otimes x$ is an idempotent of Cl_n , that is $(x \otimes x)^2 = x \otimes x$.
- (ii) if $n \equiv 0 \pmod{8}$ we have $\nu(x \otimes x) = (x \otimes x) \nu = x \otimes x$.
- (iii) for all $\varphi \in Cl_n$ we have

$$(x \otimes x) \varphi (x \otimes x) = \kappa \langle \varphi, x \otimes x \rangle (x \otimes x),$$

where $\kappa = 2^{\frac{n}{2}}$ in $n = 0 \pmod{8}$ and $\kappa = 2^{\frac{n+1}{2}}$ when $n \equiv 7 \pmod{8}$.

Proof. (i) We use (4.2) for $\psi = x$ which gives $(x \otimes x) x = x$. Therefore left multiplication of (4.2) with $x \otimes x$ gives:

$$(x \otimes x)^2 \psi = \langle \psi, x \rangle (x \otimes x) x = \langle \psi, x \rangle x = (x \otimes x) \psi,$$

for all ψ in \mathcal{S} and the claim follows.

(ii) We use the definition (4.2) to obtain after recalling that $x \in \mathcal{S}^+$

$$\nu(x \otimes x) \psi = \langle \psi, x \rangle \nu x = \langle \psi, x \rangle x = (x \otimes x) \psi,$$

for all $\psi \in \mathcal{S}$. Since $x \otimes x \in Cl_n^0$ and $n \equiv 0 \pmod{8}$, we further have $[\nu, x \otimes x] = 0$.

(iii) Let φ belong to Cl_n and $n \equiv 0 \pmod{8}$. Using again (4.2) we compute

$$(x \otimes x) \varphi (x \otimes x) \psi = \langle \psi, x \rangle (x \otimes x) \varphi x = \langle \psi, x \rangle \langle \varphi x, x \rangle x$$
$$= \langle \varphi x, x \rangle (x \otimes x) \psi$$

for all ψ in \mathcal{F} hence $(x \otimes x)\varphi(x \otimes x) = \langle \varphi x, x \rangle x \otimes x$ for all φ in Cl_n . Recall now [6] that $\langle 1, x \otimes x \rangle = 2^{-\frac{n}{2}}$ and, fact which follows essentially by taking traces and using Lemma 2.3. Therefore

$$2^{-\frac{n}{2}}\langle \varphi x, x \rangle = \langle 1, (x \otimes x)\varphi(x \otimes x) \rangle = \langle \alpha(x \otimes x)^t \cdot 1, \varphi(x \otimes x) \rangle$$
$$= \langle x \otimes x, \varphi(x \otimes x) \rangle = \langle (x \otimes x)\alpha(x \otimes x)^t, \varphi \rangle$$
$$= \langle (x \otimes x)^2, \varphi \rangle = \langle x \otimes x, \varphi \rangle$$

and the claim follows. For $n \equiv 7 \pmod{8}$ this is proved analogously.

Based on the technical Lemma above we shall compute now the holonomy algebras of the square of a spinor. Let us begin with the case of $n \equiv 0 \pmod{8}$.

Theorem 4.2. Let $n \equiv 0 \pmod{8}$ and x be a unit length spinor in S^+ and $x \otimes x$ be its square. Then:

- $\begin{array}{l} \text{(i)} \ \mathfrak{g}_{x\otimes x}^*\cong\mathfrak{so}(n,1)\\ \text{(ii)} \ Z_{x\otimes x}=\{\psi\in \mathcal{S}^+:\psi\perp x\}\oplus\{\psi\in \mathcal{S}^-:\psi\perp Vx\}. \end{array}$

Proof. (i) Let $E_T = \{X \sqcup (x \otimes x) : X \in V\}$ be the generating set of V, which is easily seen to be isomorphic to V under the map $\iota^1 : V \to E_T, \iota^1(X) = X \sqcup (x \otimes x)$. Further on, let us define $\iota^2 : \Lambda^2(V) \to Cl_n(V)$ by $\iota^2(\alpha) = \sum_{i=1}^n e_i (x \otimes x) (e_i \sqcup \alpha)$ for some orthonormal frame $\{e_i, 1 \leq i \leq n\}$ on V. This is injective since from $\iota^2(\alpha) = 0$ we get by right multiplication with $X(x \otimes x), X$ in V to $\sum_{i=1}^n e_i (x \otimes x) (e_i \sqcup \alpha) X (x \otimes x) = 0$ and further by means of Lemma 4.2 to $(X \sqcup \alpha)x \otimes x = 0$ for all X in V, whence $\alpha = 0$. Furthermore Lemma 4.2 combined with the stability relations leads easily to $[\iota^2\alpha, \iota^2\beta] = [\alpha, \beta]$ for all α, β in $\Lambda^2(V)$.

$$\begin{array}{rcl} 4\left[X \mathrel{\lrcorner} (x \otimes x), Y \mathrel{\lrcorner} (x \otimes x)\right] & = & -X(x \otimes x)Y + Y(x \otimes x)X \\ & & -(x \otimes x)XY(x \otimes x) + (x \otimes x)XY(x \otimes x) \\ & = & -\iota^2(X \wedge Y), \end{array}$$

whenever X, Y belong to V. Here we have made once more extensive use of the stability conditions and of Lemma 4.2, under the form $(x \otimes x)\Lambda^2(x \otimes x) = 0$. Hence the even commutators span $\mathfrak{so}(n)$. The triple commutator is similarly computed:

$$[\iota^{2}(\alpha), X \rfloor (x \otimes x)] = -\frac{1}{2} \sum_{i=1}^{n} \left(e_{i}(x \otimes x) F e_{i} X(x \otimes x) + (x \otimes x) X e_{i}(x \otimes x) F e_{i} - X(x \otimes x) e_{i}(x \otimes x) F e_{i} - e_{i}(x \otimes x) F e_{i}(x \otimes x) X \right)$$
$$= FX \rfloor (x \otimes x),$$

where we have made use of Lemma 4.2 and set, for convenience $\alpha = \langle F \cdot, \cdot \rangle$. Therefore $\iota^1 \oplus \iota^2 : V \oplus \Lambda^2 V \to \mathfrak{g}_{x \otimes x}^*$ gives the desired Lie algebra isomorphism between the holonomy algebra of $x \otimes x$ and $\mathfrak{so}(n, 1)$.

(ii) Let ψ belong to Z_T . Then by Lemma 3.4 this is equivalent with $(x \otimes x)\psi^+ = 0$ and $(x \otimes x)V\psi^- = 0$ and the claim follows now from the definition of $x \otimes x$, where $\psi^{\pm} \in \mathcal{S}^{\pm}$.

Remark 4.2. Squares of spinors provide examples of unipotent elements other than those coming from volume forms. Indeed, if x belongs to \mathfrak{F}^+ with |x|=1 it is easy to see that $2\sqrt{2}(x\otimes x-\frac{1}{4}(1+\nu))$ is a unipotent element in Cl_n^+ . Despite of the absence of fixed spinors in this case, the holonomy algebra remains isomorphic to $\mathfrak{so}(n,1)$.

In $n \equiv 7 \pmod{8}$ we get holonomy algebras of quite different nature than those seen before. In particular, those appear not to be perfect.

Theorem 4.3. Let $n \equiv 7 \pmod{8}$ and x belong to \mathcal{S} such that |x| = 1. Then:

- (i) $\mathfrak{g}_{x\otimes x}^*$ is abelian, isomorphic to V
- (ii) $Z_{x\otimes x}=(x)^{\perp}$.

Proof. (i) For any X, Y in V we compute

$$4[X \sqcup (x \otimes x), Y \sqcup (x \otimes x)] = (\alpha(x \otimes x)X - Xx \otimes x)(\alpha(x \otimes x)Y - Yx \otimes x)$$

$$= \left[\alpha(x \otimes x)X\alpha(x \otimes x)\right]Y - \alpha(x \otimes x)XY(x \otimes x)$$

$$- X(x \otimes x)\alpha(x \otimes x)Y + X\left[(x \otimes x)Y(x \otimes x)\right].$$

Now $x \otimes x$ belongs to Cl_n^- , hence $\alpha(x \otimes x)$ is in Cl_n^- leading to the vanishing of the second and third term above in view of the stability conditions in Lemma 2.1. Now the first and the last terms vanish too by Lemma 4.2, (iii) and since $x \otimes x$ does not contain degree 1 forms, therefore $\mathfrak{g}_{x\otimes x}^*$ is abelian. But $x\otimes x$ is non-degenerate, as it contains a non-zero multiple of the volume form whence $\mathfrak{g}_{x\otimes x}^*$ is isomorphic with V.

(ii) follows easily from the construction of
$$x \otimes x$$
.

5. 8-DIMENSIONS

In the rest of this paper we shall consider an Euclidean vector space $(V^8, \langle \cdot, \cdot \rangle)$ with orientation given by ν in $\Lambda^8(V)$. Our aim is to obtain classification results for holonomy algebras \mathfrak{g}_T^* generated by T in $\Lambda_+^4(V)$. Here we recall that in dimension 8 the Hodge start operator * preserves $\Lambda^4(V)$ which splits therefore as $\Lambda^4(V)$ = $\Lambda^4_+(V) \oplus \Lambda^4_-(V)$ into the \pm -eigenspaces of *. Our discussion is divided into several steps.

5.1. **Self-dual** 4-forms. Let us pick T in $\Lambda_+^4(V)$. Recall that in this case $Z_T = Z_T^+$ and consider the symmetric and traceless operator $\mu_T: \mathcal{S}^+ \to \mathcal{S}^+$. Let $\sigma_T = \{\lambda_q, 1 \leq$ $q \leq p$ be the non-zero part of the spectrum of μ_T where we assume the eigenvalues $\lambda_q, 1 \leq q \leq p$ to be pairwise distinct and where we denote their multiplicities by $m_q, 1 \leq q \leq p$. Therefore we obtain a splitting

$$(5.1) \mathcal{S}^+ = Z_T \oplus \mathcal{S}_1 \oplus \dots \mathcal{S}_p$$

where \mathcal{S}_q are the eigenspaces of μ_T corresponding to the eigenvalues $\lambda_q, 1 \leq q \leq p$. Our aim here is to examine the splitting of Λ^2 induced by (5.1) and to relate it directly to the form T. We need now to recall the following simple fact, which essentially exploits the squaring isomorphism in 8-dimensions.

Lemma 5.1. Let x, y belong to S^+ , and let $x \wedge y$ in Cl_8 be given as

$$(x \wedge y)\psi = \langle \psi, x \rangle y - \langle \psi, y \rangle x$$

for all ψ in \mathcal{S} . Then:

- (i) $x \wedge y$ belongs to $Cl_8^0 \cap Cl_8^+$ and $(x \wedge y)^t = -x \wedge y$. (ii) if moreover $Tx = \lambda_1 x$ and $Ty = \lambda_2 y$ where T belongs to Λ_+^4 then

$$T(x \wedge y)T = \lambda_1 \lambda_2 x \wedge y$$
, and $T(x \wedge y) + (x \wedge y)T = (\lambda_1 + \lambda_2)x \wedge y$.

(iii) under the assumptions in (ii), if $\lambda_1 = \lambda_2$ then $T(x \wedge y) = (x \wedge y)T = \lambda_1 x \wedge y$.

(iv) if x', y' is another pair of spinors in \mathcal{S}^+ then

$$[x \wedge y, x' \wedge y'] = \langle x, y' \rangle x' \wedge y - \langle y, y' \rangle x' \wedge x - \langle x, x' \rangle y' \wedge y + \langle x', y \rangle y' \wedge x.$$

(v)
$$8\langle x \wedge y, x' \wedge y' \rangle = \langle y, y' \rangle \langle x, x' \rangle - \langle x, y' \rangle \langle y, x' \rangle.$$

Proof. (i) is standard, see [13].

We prove (ii) and (iii) at the same time. For any ψ in \mathcal{S} we have

$$(x \wedge y)T\psi = \langle T\psi, x \rangle y - \langle T\psi, y \rangle x = \langle T\psi, x \rangle y - \langle T\psi, y \rangle x = \lambda_1 \langle \psi, x \rangle y - \lambda_2 \langle \psi, y \rangle x$$
as $\langle T\psi, x \rangle = \langle \psi, Tx \rangle = \lambda_1 \langle \psi, x \rangle$ and similarly $\langle T\psi, y \rangle = \lambda_2 \langle \psi, y \rangle$. Moreover,

$$T(x \wedge y)\psi = \langle \psi, x \rangle Ty - \langle \psi, y \rangle Tx = \lambda_2 \langle \psi, x \rangle y - \lambda_1 \langle \psi, y \rangle x.$$

All claims in (ii) and (iii) follow now easily. The proof of (iv) is a straightforward direct computation involving only the definition of the exterior product of spinors. (v) A direct computation based on the definition of the wedge product of spinors shows that the trace of the Clifford multiplication with the symmetric (in the sense of Lemma 2.3) element $(x \wedge y)(x' \wedge y') + (x' \wedge y')(x \wedge y)$ of Cl_8 is given by

$$-4\Big[\langle y,y'\rangle\langle x,x'\rangle-\langle x,y'\rangle\langle y,x'\rangle\Big].$$

The claim follows now by using Lemma 2.3.

For any $1 \le k, i, j \le p$ let us now define the spaces

$$E_k = \{ \alpha \in \Lambda^2 : T\alpha T = 0, \ T\alpha + \alpha T = \frac{\lambda_k}{2} (1 + \nu)\alpha \}$$

$$F_{ij} = \{ \alpha \in \Lambda^2 : T\alpha T = \frac{1}{2} \lambda_i \lambda_j (1 + \nu)\alpha, \ T\alpha + \alpha T = \frac{\lambda_i + \lambda_j}{2} (1 + \nu)\alpha \}$$

$$\iota_T^0 = \{ \alpha \in \Lambda^2 : \alpha T = T\alpha = 0 \}.$$

Obviously we have $F_{ij} = F_{ji}$. For notational convenience, we set $E = \bigoplus_{k=1}^{p} E_k$, $F = \bigoplus_{k=1}^{p} E_k$

 $\bigoplus_{1 \le i \le j \le p} F_{ij}$. Another related object is

Definition 5.1. The isotropy algebra ι_T of T in Λ^4 is the subalgebra of $\mathfrak{so}(V)$ given by

$$\{\alpha\in\mathfrak{so}(V): [\alpha,T]=0\}.$$

Here the Lie bracket is considered within the Lie algebra Cl₈.

Proposition 5.1. The following hold:

(i) there is an orthogonal, direct sum decomposition

$$\Lambda^2=\iota^0_T\oplus E\oplus F$$

(ii) we have the following string of isomorphisms:

$$\iota_T^0 \cong \Lambda^2(Z_T)
E_k \cong Z_T \otimes \mathcal{S}_k
F_{kk} \cong \Lambda^2(\mathcal{S}_k)
F_{ij} \cong \mathcal{S}_i \otimes \mathcal{S}_j, i \neq j$$

(iii)
$$\iota_T = \iota_T^0 \oplus \bigoplus_{k=1}^p F_{kk}$$

Proof. We prove (i) and (ii) together. For any two vector sub-spaces V, W of \mathcal{S}^+ we denote by $V \hat{\otimes} W$ the inclusion of $V \otimes W$ into $\Lambda^2(\mathcal{S}^+)$. Letting now ι_T^0, E_k, F_{ij} be the images of $\Lambda^2(Z_T)$, $Z_T \hat{\otimes} \mathcal{S}_k$, $\mathcal{S}_i \hat{\otimes} \mathcal{S}_j$ under the inverse of the linear isomorphism $\mu_{|(1+\nu)\Lambda^2}: (1+\nu)\Lambda^2 \to \Lambda^2(\mathcal{S}^+)$ proves our claims by making use of Lemma 5.1. (iii) Pick α in ι_T . Then $T\alpha = \alpha T$ hence $T^2\alpha + \alpha T^2 = 2T\alpha T$. It is easy to see that the operator $\{T^2,\cdot\} - 2T \cdot T$ equals 0 on $\iota_T^0, \lambda_k^2 1_{E_k}$ on E_k and $(\lambda_i - \lambda_j)^2 1_{E_{ij}}$ on E_{ij} thus $\iota_T \subseteq \iota_T^0 \oplus \bigoplus_{k=1}^p F_{kk}$. The reverse inclusion and therefore the equality follows from the construction of the spaces F_{kk} , $1 \le k \le p$ and Lemma 5.1, (iii).

In the Proposition above the fact that V is 8-dimensional, which implies that $A^0 \cap Cl_8^+ = (1+\nu)\Lambda^2$ had been used in a crucial way. It is obviously valid on Λ^- as well. The block structure of the isotropy algebra of a form T in Λ_+^4 has been already obtained in [5], by a slightly different method and under the additional assumption that T is a calibration on V. In this case, the work in [5] gives a complete geometric description of the resulting orbits. In order to understand the structure of the holonomy algebra of T we need to have a look at the Lie algebraic features of the splitting above.

Corollary 5.1. Let T in $\Lambda_+^4(V)$ with $Z_T \neq (0)$ be given. Then F is a Lie sub-algebra of Λ^2 isomorphic with $\mathfrak{so}(Z_T^{\perp})$.

Proof. From the construction of F the Clifford multiplication map gives an isometry $\mu: (1+\nu)F \to \Lambda^2(Z_T^{\perp})$ by Lemma 5.1, (v). Moreover, this is a Lie algebra isomorphism by (iv) of the same Lemma.

Lemma 5.2. Suppose that $Z_T \neq (0)$. We have :

- (i) $[\iota_{T}^{0}, \iota_{T}^{0}] = \iota_{T}^{0}, \ [\iota_{T}^{0}, E_{k}] = E_{k}, \ [\iota_{T}^{0}, F_{ij}] = 0 \ for \ any \ 1 \leq k, i, j \leq p.$ (ii) $[E_{i}, E_{j}] = F_{ij} \ for \ i \neq j \ and \ [E_{i}, E_{i}] = F_{ii} \oplus \iota_{T}^{0}$ (iii) $[E_{i}, F_{jk}] = 0 \ if \ i \neq j, k \ and \ [E_{i}, F_{ij}] = E_{j}$

- (iv) $[F_{ij}, F_{ks}] = 0$ if $(i, j) \cap (k, s) = \emptyset$
- (v) $[F_{ij}, F_{ik}] = F_{jk}$ if i, j, k are mutually distinct (vi) $[F_{ij}, F_{ij}] = F_{ii} \oplus F_{jj}$ and $[F_{ii}, F_{ik}] = F_{ik}$ provided that $m_i \geq 2$.

Proof. Follows directly from the general formula in (iii) of Lemma 5.1 after inspecting the various possibilities of combining factors in the splitting of \mathcal{S}^+ as given by

We need now to establish the analogue of Lemma 5.2 when the form T has $Z_T = (0)$. In this case $\iota_T^0 = E = (0)$ and using the same arguments as previously

Lemma 5.3. Let T be in Λ_+^4 with $Z_T = (0)$. We have :

- (i) $[F_{ij}, F_{ks}] = 0$ if $(i, j) \cap (k, s) = \emptyset$
- (ii) $[F_{ij}, F_{ik}] = F_{jk}$ if i, j, k are mutually distinct (iii) $[F_{ij}, F_{ij}] = F_{ii} \oplus F_{jj}$ and $[F_{ii}, F_{ik}] = F_{ik}$ provided that $m_i \geq 2$.

5.2. Structure of the commutators. We shall give in this section a simplified expression, relying on the particular dimension, for the generating space of the even part of the holonomy algebra \mathfrak{g}_T^* , where T belongs to Λ_+^4 .

Lemma 5.4. Let T belong to $Cl_8^0 \cap Cl_8^+$ satisfy $T^t = T$. We have

$$-4[X \rfloor T, Y \rfloor T] = 2T\alpha T + \frac{1}{4}L(T^{2}\alpha + \alpha T^{2}) + 4|T|^{2}(1-\nu)\alpha$$

for all X, Y in V, where α in $\Lambda^2(V)$ is given as $\alpha = X \wedge Y$.

Proof. Let $a: \Lambda^2(V) \to Cl_8$ be defined by setting

$$a(\alpha) = \sum_{i=1}^{8} e_i T^2(e_i \, \lrcorner \, \alpha)$$

for some orthonormal frame $\{e_i, 1 \leq i \leq 8\}$. Lemma 4.1 actually says that

$$-4[X \rfloor T, Y \rfloor T] = 2T\alpha T + a(\alpha)$$

and we need only work out a simpler expression for the operator a. We compute

$$e_i T^2(e_i \perp \alpha) = \frac{1}{2} e_i (T^2 \alpha) e_i - \frac{1}{2} (e_i T^2 e_i) \alpha$$

leading to $a(\alpha) = \frac{1}{2}L(T^2\alpha) - \frac{1}{2}L(T^2)\alpha$. But $L(T^2\alpha) = \frac{1}{2}L(T^2\alpha + \alpha T^2)$ since $[T^2, \alpha]$ is a 4-form, and moreover $LT^2 = -8|T|^2(1-\nu)$.

5.3. Computation of $\mathfrak{g}_{T}^{*,0}$. For a given T in $\Lambda_{+}^{4}(\mathbb{R}^{8})$, we shall compute now the even part $\mathfrak{g}_{T}^{*,0}$ of its holonomy algebra. The main technical ingredient in this section is contained in the following observation.

Lemma 5.5. Let T be in Λ_+^4 . Then $\mathfrak{g}_{T^{2k+1}}^* \subseteq \mathfrak{g}_T^*$ for all k in \mathbb{N} .

Proof. Let $\{e_i, 1 \leq i \leq 8\}$ be an orthonormal basis in V and consider the partial Casimir operator $C_T : Cl_8 \to Cl_8$ given by

$$C_T = \sum_{i=1}^{8} [e_i \, \lrcorner \, T, [e_i \, \lrcorner \, T, \cdot]].$$

Obviously, C_T preserves the algebra \mathfrak{g}_T^* , that is $C_T(\mathfrak{g}_T^*) \subseteq \mathfrak{g}_T^*$. A straightforward computation actually shows that

$$C_T \varphi = \sum_{i=1}^8 (e_i \, \rfloor \, T)^2 \varphi + \varphi(e_i \, \rfloor \, T)^2 - 2(e_i \, \rfloor \, T) \varphi(e_i \, \rfloor \, T)$$

for all φ in Cl_8 . We shall now compute $C_T(X \perp \varphi)$ where φ belongs to $Cl_8^0 \cap Cl_8^+$ is such that $\varphi^t = \varphi$ (equivalently φ is in $\Lambda_+^4 \oplus \mathbb{R}(1 + \nu)$). We compute

$$-8(e_i \sqcup T)(X \sqcup \varphi)(e_i \sqcup T) = (e_i T - Te_i)(X\varphi - \varphi X)(e_i T - Te_i)$$

$$= (-Te_i X\varphi - e_i T\varphi X)(e_i T - Te_i)$$

$$= -e_i T\varphi X e_i T + Te_i X\varphi T e_i$$

$$= -e_i T\varphi (-2\langle e_i, X \rangle - e_i X)T + T(-2\langle e_i, X \rangle - Xe_i)\varphi Te_i$$

henceforth after summation we get

$$4\sum_{i=1}^{8} (e_i \rfloor T)(X \rfloor \varphi)(e_i \rfloor T) = -[X, T\varphi T] - \frac{1}{2}L(T\varphi)XT + \frac{1}{2}TXL(\varphi T).$$

Now

$$4\sum_{i=1}^{8} (e_i \perp T)^2 = \sum_{i=1}^{8} (e_i T - Te_i)^2 = LTT - LT^2 + 8T^2 + TLT = 8T^2 + 8|T|^2(1 - \nu)$$

as LT = 0 and $LT^2 = -8|T|^2(1 - \nu)$. A short computation using the stability relations gives now

$$T^2(X \mathrel{\lrcorner} \varphi) + (X \mathrel{\lrcorner} \varphi)T^2 = -\tfrac{1}{2}(-T^2\varphi X + X\varphi T^2)$$

hence in the end we obtain

$$4C_T(X \perp \varphi) = -4(X\varphi T^2 - T^2\varphi X) + 2[X, T\varphi T] + L(T\varphi)XT - TXL(\varphi T) + 16|T|^2(X \perp \varphi)$$

for all X in V and where φ belongs to $\Lambda_+^4 \oplus \mathbb{R}(1+\nu)$. In particular, for $\varphi = T^k, k$ in \mathbb{N} this yields

$$4C_T(X \perp T^k) = X \perp T^{k+2} + 32\langle T^k, T \rangle X \perp T + 16|T|^2(X \perp T^k)$$

for all X in V, where we have use that $L(T^{k+1}) = -8\langle T^k, T \rangle (1-\nu)$. By induction, given that C_T preserves \mathfrak{g}_T^* and that the later contains $\{X \, \rfloor \, T : X \in V\}$ we arrive at $\{X \, \rfloor \, T^{2k+1} : X \in V\} \subseteq \mathfrak{g}_T^*$, for all k in $\mathbb N$ and our claim follows. \square

Lemma 5.6. Let λ_i where $1 \leq i \leq p$ belong to σ_T . If $\dim_{\mathbb{R}} Z_T \neq 0$ and

$$8|T|^{2k+1} = \lambda_i^{2(2k+1)}$$

holds for all k in \mathbb{N} then $dim_{\mathbb{R}}Z_T = 6$ and $\sigma_T = \{\lambda_i, -\lambda_i\}$ with multiplicities (1, 1), provided that $T \neq 0$.

Proof. Since $16|T|^{2k+1} = \sum_{q=1}^{p} m_q \lambda_q^{2(2k+1)}$ by making use of Lemma 2.3, the equation we have to solve becomes

(5.2)
$$\sum_{q=1}^{p} m_q \lambda_q^{2(2k+1)} = 2\lambda_i^{2(2k+1)}$$

for all k in \mathbb{N} . We now divide by $\lambda_i^{2(2k+1)}$ and take the limit when $k \to \infty$. It follows that $|\lambda_q| \le |\lambda_i|$ for all $1 \le q \le i$ and also that $\sum_{|\lambda_q| = |\lambda_i|} m_q = 2$. It follows easily that

 $m_i = 1$, otherwise we would have $m_i = 2$ and further $\sigma_T = \{\lambda_i\}$ by making use of (5.2), which contradicts that μ_T is traceless. Therefore $-\lambda_i$ belongs to σ_T , with multiplicity 1 and our claim follows again from (5.2).

Proposition 5.2. Let T belong to $\Lambda^4_+(V)$ with $dim_{\mathbb{R}}Z_T \neq 0, 6$. Then

$$(1+\nu)F \oplus (1-\nu)\Lambda^2 \subseteq \mathfrak{g}_T^{*,0}$$
.

Proof. Making use of Lemma 5.4 we have that $\mathfrak{g}_T^{*,0}$ contains the set

(5.3)
$$\{2T\alpha T + \frac{1}{4}L(T^2\alpha + \alpha T^2) + 4|T|^2(1-\nu)\alpha : \alpha \in \Lambda^2\}.$$

as this is just spanned by double commutators of elements in its generating set. From the above we find that $(1-\nu)\iota^0_T$ is contained in $\mathfrak{g}^{*,0}_T$. Actually, by using Lemma 5.5 we have that $X \, \lrcorner \, T^{2k+1}$ belongs to \mathfrak{g}^*_T and therefore, after taking double commutators of such elements and using again Lemma 5.4 we get that

$$2T^{2k+1}\alpha T^{2k+1} + \frac{1}{4}L(T^{2(2k+1)}\alpha + \alpha T^{2(2k+1)}) + 4|T^{2k+1}|^2(1-\nu)\alpha$$

belongs to $\mathfrak{g}_T^{*,0}$ for any α in Λ^2 . Now if α is in E_i , for some $1 \leq i \leq p$ we have $T\alpha T = 0$ and an easy computation by induction shows

$$T^{2(2k+1)}\alpha + \alpha T^{2(2k+1)} = \frac{1}{2}\lambda_i^{2(2k+1)}(1+\nu)\alpha$$

for all k in \mathbb{N} . We are led eventually to having

$$(4|T^{2k+1}|^2 - \frac{1}{2}\lambda_i^{2(2k+1)})(1-\nu)E_i$$

contained in \mathfrak{g}_T^* for all $1 \leq i \leq p$ and all k in \mathbb{N} .

Now since σ_T has not the form in Lemma 5.6, in other words $dim_{\mathbb{R}}Z_T \neq 6$, for each $1 \leq i \leq p$ the factor above will be non-vanishing for some k in N whence $(1-\nu)E_i \subseteq \mathfrak{g}_T^*$ whenever $1 \leq i \leq p$. Now taking commutators and using (ii) of Lemma 5.2 it follows that $(1-\nu)\Lambda^2 \subseteq \mathfrak{g}_T^*$. But $L(T^2\alpha + \alpha T^2)$ belongs to $(1-\nu)\Lambda^2$ for all α in Λ^2 hence we get from (5.3) that $\{T\alpha T : \alpha \in \Lambda^2\}$ is contained in $\mathfrak{g}_T^{*,0}$. Making use of the splitting in Proposition 5.1 this actually says that $(1+\nu)F\subseteq \mathfrak{g}_T^*$ and we have showed that

$$(1+\nu)F \oplus (1-\nu)\Lambda^2 \subseteq \mathfrak{g}_T^{*,0}.$$

Therefore, when the set of spinors fixed by some self-dual 4-form is not empty we can conclude, with one exception, that:

Proposition 5.3. Let T in $\Lambda^4_+(V)$ be given, and suppose that $\dim_{\mathbb{R}} Z_T \neq 0, 6$. Then

(i)
$$\mathfrak{g}_{T}^{*,0} = (1+\nu)F \oplus (1-\nu)\Lambda^{2}$$

(ii) $Z(\mathfrak{g}_{T}^{*,0}) = (0)$.

(ii)
$$Z(\mathfrak{g}_T^{*,0}) = (0).$$

Proof. (i) By Proposition 5.2, it is enough to see that $\mathfrak{g}_T^{*,0} \subseteq (1+\nu)F \oplus (1-\nu)\Lambda^2$ and this will be achieved by showing that $\mathfrak{g}_T^{*,0}$ is orthogonal to $(1+\nu)(\iota_0^T \oplus E)$. Indeed, by the definition of Z_T we have $(X \perp T)Z_T = 0$ for all X in V, therefore $\mathfrak{g}_T^* = 0$. We now pick φ in $\mathfrak{g}_T^{*,0}, x$ in Z_T and y in \mathcal{F}^+ . From the definition of $x \wedge y$ and $\varphi x = 0$ follows that

$$\varphi(x \wedge y)\psi = \langle x, \psi \rangle \varphi y$$
$$(x \wedge y)\varphi \psi = -\langle \varphi \psi, y \rangle x = \langle \psi, \varphi y \rangle x$$

for all ψ in \mathcal{S} . Since $\varphi(x \wedge y) + \varphi(x \wedge y)$ is a symmetric element in Cl_8 after taking the trace we get

$$Tr(\varphi(x \wedge y) + (x \wedge y)\varphi) = 2\langle \varphi y, x \rangle = -2\langle y, \varphi x \rangle = 0.$$

Now using Lemma 2.3 it follows that $\langle \varphi(x \wedge y) + (x \wedge y) \varphi, 1 \rangle = 0$ and since $\alpha(\varphi^t) = -\varphi$ we arrive at $\langle \varphi, x \wedge y \rangle = 0$. Because $\{x \wedge y : x \in Z_T, y \in \mathcal{F} \text{ spans } (1 + \nu)(\iota_T^0 \oplus E) \text{ it follows that } \mathfrak{g}_T^{*,0} \text{ is orthogonal to } (1 + \nu)(\iota_T^0 \oplus E), \text{ hence contained in } (1 + \nu)F \text{ and the claim follows.}$

We leave out now the case when $dim_{\mathbb{R}}Z_T=6$, to be treated further on, and look at the situation when there are no non-zero fixed spinors which needs first a combinatorics Lemma.

Lemma 5.7. Let λ_i, λ_j belong to σ_T with $1 \leq i \neq j \leq p$. If $\dim_{\mathbb{R}} Z_T = 0$ and

$$4|T^{2k+1}|^2 - \frac{1}{2}(\lambda_i^{2(2k+1)} + \lambda_j^{2(2k+1)}) = (\lambda_i \lambda_j)^{2k} (4|T|^2 - \frac{1}{2}(\lambda_i^2 + \lambda_j^2))$$

holds for all k in \mathbb{N} then σ_T must be one of the following:

- (i) $\sigma_T = \{\lambda, -\lambda\}$
- (ii) $\sigma_T = \{\lambda, -\lambda, \mu, -\mu, \sqrt{|\lambda\mu|}, -\sqrt{|\lambda\mu|}\}$ with multiplicities (1, 1, 1, 1, 2, 2), for some $|\lambda| \neq |\mu|$
- (iii) $\sigma_T = \{\lambda, -\mu, \mu, \pm \sqrt{|\lambda \mu|}\}\$ with multiplicities (2, 1, 1, 4).

Proof. The proof is given in the Appendix for it is of rather technical nature and the reader may skip it at a first reading. \Box

Proposition 5.4. Let T belong to Λ_+^4 such that $Z_T = (0)$. Then either

(i)

$$\mathfrak{g}_T^{*,0} = A^0 \cong \mathfrak{so}(8) \oplus \mathfrak{so}(8)$$

or

(ii)

$$\mathfrak{g}_T^{*,0} \cong \mathfrak{so}(8)$$

case which occurs when T is a unipotent element of Cl_8^+ , in the sense that $T^2 = \lambda(1+\nu)$ for some $\lambda > 0$.

Proof. Using Lemma 5.5 we have that $[X \rfloor T^{2k+1}, Y \rfloor T^{2k+1}]$ belongs to $\mathfrak{g}_T^{*,0}$ for all X, Y in V and any natural number k. Now Lemma 5.4 implies that

$$2T^{2k+1}\alpha T^{2k+1} + \frac{1}{4}L(T^{2(2k+1)}\alpha + \alpha T^{2(2k+1)}) + 4|T^{2k+1}|^2(1-\nu)\alpha$$

belongs to \mathfrak{g}_T^* for any α in $\Lambda^2(V)$ and any k in \mathbb{N} . Let now $1 \leq i, j \leq p$ and pick α in F_{ij} . An easy computation by induction gives

$$T^{2k+1}\alpha T^{2k+1} = \frac{1}{2}(\lambda_i \lambda_j)^{2k+1} (1+\nu)\alpha$$

$$T^{2(2k+1)}\alpha + \alpha T^{2(2k+1)} = \frac{1}{2}(\lambda_i^{2(2k+1)} + \lambda_j^{2(2k+1)})(1+\nu)\alpha$$

hence

(5.4)
$$(\lambda_i \lambda_j)^{2k+1} (1+\nu)\alpha + \left[4|T^{2k+1}|^2 - \frac{1}{2}(\lambda_i^{2(2k+1)} + \lambda_j^{2(2k+1)})\right] (1-\nu)\alpha$$

belongs to $\mathfrak{g}_T^{*,0}$ for all α in F_{ij} and all k in \mathbb{N} . But in the same time, for k=0, we have

$$\lambda_i \lambda_j (1+\nu)\alpha + \left[4|T|^2 - \frac{1}{2}(\lambda_i^2 + \lambda_j^2)\right](1-\nu)\alpha$$

belongs to $\mathfrak{g}_T^{*,0}$ for all α in F_{ij} hence

$$\left[4|T^{2k+1}|^2 - \frac{1}{2}(\lambda_i^{2(2k+1)} + \lambda_j^{2(2k+1)}) - (\lambda_i \lambda_j)^{2k}(4|T|^2 - \frac{1}{2}(\lambda_i^2 + \lambda_j^2))\right](1 - \nu)\alpha$$

is in $\mathfrak{g}_T^{*,0}$ whenever α belongs to F_{ij} and k is in \mathbb{N} . If T is not an unipotent element, nor σ_T has the form in (ii) or (iii) of Lemma 5.7 we use Lemma 5.7 to see that the factor above will be non vanishing for some k. This leads to $(1-\nu)F_{ij} \subseteq \mathfrak{g}_T^{*,0}$ and further from (5.4) we get that $(1+\nu)F_{ij} \subseteq \mathfrak{g}_T^{*,0}$, since $0 \notin \sigma_T$. Therefore when the spectrum of μ_T is not as in (ii) or (ii) of Lemma 5.7 our claim follows easily from the above, while the case when T is unipotent is covered by Theorem 4.1.

It remains now to treat the remaining two cases both of which to be worked out directly from (5.4).

(a) $\sigma_T = \{\lambda, -\lambda, \mu, -\mu, \sqrt{|\lambda\mu|}, -\sqrt{|\lambda\mu|}\}$ with multiplicities (1, 1, 1, 1, 2, 2). In this case we have

$$4|T|^{2k+1} = \frac{1}{2}(\lambda^{2(2k+1)} + \mu^{2(2k+1)}) + |\lambda\mu|^{2k+1}$$

for all k in \mathbb{N} . We shall label the distinct eigenvalues as $\lambda_i, 1 \leq i \leq 6$ in the order they are listed in σ_T . For any α in $F_{i5}, 1 \leq i \leq 6$ we have then from (5.4)

$$2(\lambda_i \sqrt{|\lambda \mu|})^{2k+1} (1+\nu)\alpha + \left[(\lambda^{2(2k+1)} + \mu^{2(2k+1)}) + |\lambda \mu|^{2(2k+1)} - \lambda_i^{2(2k+1)} \right] (1-\nu)\alpha$$

for all k in \mathbb{N} . w.l.o.g we may also assume that $|\lambda| < |\mu|$. Now if i = 1, 2 or i = 5, 6 dividing by $\mu^{2(2k+1)}$ and taking the limit with $k \to \infty$ yields $(1-\nu)F_{i5} \subseteq \mathfrak{g}_T^{*,0}$ whence the same holds for $(1+\nu)F_{i5}$. Now if i = 3, 4 we divide by $(\mu\sqrt{|\lambda\mu|})^{2k+1}$ and make $k \to \infty$ to arrive at $(1+\nu)F_{i5}$ contained in $\mathfrak{g}_T^{*,0}$ and from there to $(1-\nu)F_{i5} \subseteq \mathfrak{g}_T^0$. Now taking commutators and using Lemma 5.3 it follows that $(1\pm\nu)F_{ij}$ is contained in $\mathfrak{g}_T^{*,0}$ hence the latter contains A^0 and our claim follows.

(b) $\sigma_T = \{\lambda, -\mu, \mu, \pm \sqrt{|\lambda \mu|}\}\$ with multiplicities (2, 1, 1, 4).

Here we label the (distinct) eigenvalues of μ_T by $\lambda_i, 1 \leq i \leq 5$, in their order of appearance and note that from the traceless of μ_T we must have $|\lambda| = 4|\mu| > |\mu|$. Apart from this differences the proof of (a) continues to hold without any change. We have exhausted all possibilities and therefore our claim is finally proved. \square

We conclude this section with giving the full description of the holonomy algebras of forms T in Λ_+^4 with $Z_T = (0)$.

Theorem 5.1. Let T in Λ_+^4 be given and suppose moreover that $Z_T = (0)$. Then either

$$\mathfrak{g}_T^* = A \cong \mathfrak{so}(8,8)$$

or

$$\mathfrak{g}_T^* \cong \mathfrak{so}(8,1).$$

The latter case occurs when T is an unipotent element of Cl_8^+ , that is $T^2 = \lambda(1+\nu)$ for some $\lambda > 0$.

Proof. Because $[\mathfrak{g}_T^{*,0},\mathfrak{g}_T^{*,1}]\subseteq \mathfrak{g}_T^{*,1}$, when $\mathfrak{g}_T^{*,0}=A^0$ we obtain that $\mathfrak{g}_T^{*,1}$ is an invariant sub-space of A^1 , w.r.t the adjoint representation of A^0 on A^1 . Because this is irreducible (see Lemma 3.2, (iii)) we find that $\mathfrak{g}_T^{*,1}=A^1$ whence $\mathfrak{g}_T^*=A$. The case of an unipotent element has been treated in Theorem 4.1.

5.4. The full holonomy algebra when $Z_T \neq (0)$. In order to have a complete description of holonomy algebras of self-dual 4-forms in 8-dimensions it remains to understand the odd part $\mathfrak{g}_T^{*,1}$ of \mathfrak{g}_T^* for some T in Λ_+^4 when $Z_T \neq (0)$. We recall that in this situation $\mathfrak{g}_T^{*,0}$ has been computed in Proposition 5.3. Let us now define

$$Q = \{ \varphi \in A^1 : T\varphi + \varphi T = 0 \}.$$

We shall also work with the symmetric tensor product of spinors $(x, y) \to x \odot y$ where x, y belongs to \mathcal{S}^+ which is defined in analogy with the exterior product of spinors we saw before.

Lemma 5.8. Let T belong to Λ^4_+ . The following hold:

- (i) $Q = \{ \varphi \in A : T\varphi = \varphi T = 0 \}.$
- (ii) the map $(x,y) \to x \odot y$ extends to an isomorphism $\mathcal{G}^- \otimes Z_T \to Q$.

Proof. (i) If $T\varphi + \varphi T = 0$ with φ in A^1 , left multiplication with ν gives $T\varphi - \varphi T = 0$, hence our claim, while using that $\nu T = T$ and $\nu \varphi + \varphi \nu = 0$.

(ii) It is easy to see from (i) that for any φ in Q the map μ_{φ} is a symmetric endomorphism of \mathcal{S} such that $\mu_{\varphi}\mathcal{S}^- \subseteq Z_T$. Details are very similar to previous proofs and therefore left to the reader.

Proposition 5.5. Let T belong to Λ_+^4 with $Z_T \neq (0)$. We have that $\mathfrak{g}_T^{*,1} = Q^{\perp}$.

Proof. A direct computation shows that

$$-[(1-\nu)\beta,X \,\lrcorner\, T] = (\alpha \wedge X)T + T(\alpha \wedge X) + [X \,\lrcorner\, \alpha,T]$$

for all α, β in Λ^2 and all X in V. Given that $(1 - \nu)\Lambda^2 \subseteq \mathfrak{g}_T^{*,0}$ it follows easily that $T\varphi + \varphi T$ belongs to $\mathfrak{g}_T^{*,1}$ for all φ in Λ^3 and further that this actually holds for all φ in $A^1 = \Lambda^3 \oplus \Lambda^7$. This is because $T\Lambda^7 + \Lambda^7 T$ just gives the generating set of \mathfrak{g}_T^* since $\Lambda^7 = \nu \Lambda^1$.

Therefore $\mathfrak{g}_T^{*,1}$ contains the image of the symmetric operator $\{T,\cdot\}:A^1\to A^1$ hence Q^\perp . Using Lemma 5.8, an argument completely similar to the one in the proof of Proposition 5.3, (i) leads to $Q^\perp\subseteq\mathfrak{g}_T^{*,1}$ hence to the proof of the claim.

Therefore, our main in result in this section, our main result on holonomy algebras of self-dual 4-forms with fixed spinors from this section is

Theorem 5.2. Let T be in Λ_+^4 with $dim_{\mathbb{R}}Z_T \neq 0, 6$. Then the Clifford multiplication realises a Lie algebra isomorphism

$$\mu: \mathfrak{g}_T^* \to \mathfrak{so}(8, 8 - dim_{\mathbb{R}} Z_T).$$

Proof. It is enough to see that the map $(x,y) \to x \odot y$ gives an isomorphism $\mathcal{F}^- \otimes Z_T^\perp \to Q^\perp$ and the rest follows by collecting the results above.

It is also easy to see that under the assumptions above, \mathfrak{g}_{T}^{*} is a perfect Lie algebra.

6. The case when $dim_{\mathbb{R}}Z_T=6$

In this section we shall continue to work on an 8-dimensional Euclidean vector space $(V^8, \langle \cdot, \cdot \rangle)$ which is furthermore supposed to be oriented, with orientation form given by ν in $\Lambda^8(V)$. Moreover, we will assume that T in $\Lambda_+^4(V)$ satisfies $\dim_{\mathbb{R}} Z_T = 6$, and our primary aim will be to compute, in this case, the algebra \mathfrak{g}_T^* , the only case left out in the previous sections. Since this situation cannot be covered only by the methods there we need to develop some more material concerning the structure of such forms. Let therefore $\sigma_T = \{\lambda_1, \lambda_2\}$ be the non-zero part of the spectrum of $\mu_T : \mathcal{F}^+ \to \mathcal{F}^+$ with multiplicities (1,1) and let us also recall that $\lambda_1 + \lambda_2 = 0$. We equally recall that in this case the splitting of Λ^2 from Proposition 5.1 becomes

(6.1)
$$\Lambda^2 = \iota_T^0 \oplus E_1 \oplus E_1 \oplus F_{12}$$

and in particular F is reduced to the 1-dimensional component F_{12} . In what follows we shall use the normalisation $\lambda_1 = 1$ as it is clear that rescaling the generating form leaves a holonomy algebra unchanged.

6.1. **Spinor** 2-planes. We start by recalling the following

Definition 6.1. Let $(V^8, \langle \cdot, \cdot \rangle)$ be an Euclidean vector space. An almost Hermitian structure consists in a linear almost complex structure J which is orthogonal w.r.t. the scalar product $\langle \cdot, \cdot \rangle$. If moreover V is oriented, with orientation given by ν in $\Lambda^n(V)$, J is positive if $\omega^4 = \lambda \nu$ for some $\lambda > 0$ where $\omega = g(J \cdot, \cdot)$ is the so-called Kähler form of (g, J).

In what follows we shall keep all previous notations and also recall the following well known fact, see [6] for instance.

Proposition 6.1. Let $L \subseteq \mathcal{S}^+$ be any oriented 2-dimensional sub-space of spinors. Then L determines a unique positive almost Hermitian structure, say J, on V.

For later use, and by sending again the reader to [6] $\{x_1, x_2\}$ in L with the convention that $|x_1|^2 = |x_2|^2 = 2$, where $\omega = \langle J \cdot, \cdot \rangle$. It is not difficult to see that the converse of Proposition 6.1 also holds, in the sense that any compatible almost Hermitian structure J defines a 2-dimensional sub-space L of \mathcal{F}^+ which is explicitly given as $L = Ker(\mu_\omega^2 + 16)$. For any compatible almost Hermitian structure J we denote by λ^4 the underlying real bundle of the canonical line bundle of J. Explicitly, $\lambda^4 = \{\alpha \in \Lambda^4 : \alpha(J \cdot, J \cdot, \cdot, \cdot) = -\alpha\}$ and if moreover J is positive λ^4 is contained in Λ^+ (see [14]). Note that if the contrary is not specified all forms are real valued in this setting. We also recall that in presence of an almost Hermitian structure $\Lambda^2 = \Lambda_0^2 \oplus \mathbb{R}\omega$, an orthogonal, direct, sum and that $\lambda^{1,1}$ is the space of J-invariant 2-forms on V.

Lemma 6.1. Let $L \subseteq \mathcal{S}^+$ be two dimensional and oriented and let J be the complex structure determined by L. Then T in $\Lambda^2(V)$ satisfies TL = 0 iff T belongs to $\lambda_0^{1,1}(V)$.

Proof. This is an easy exercise taking into account that from the construction of Jit follows

$$JYx_1 = -Yx_2$$
$$JYx_2 = Yx_1$$

for all Y in V, where $\{x_1, x_2\}$ is an oriented orthonormal basis in L.

This essentially leads to having $\iota_T^0 = \lambda_0^{1,1}$ fact to be used later on and which encodes the well-known special isomorphism $\mathfrak{su}(4) \cong \mathfrak{so}(6)$ [13]. Moving within the same circle of arguments also shows that

Proposition 6.2. Given any 2-dimensional sub-space $L \subseteq \mathcal{S}^+$, the map $(x,y) \to$ $x \odot y$ extends to an isomorphism $S_0^2(L) \to \lambda^4$.

Proposition 6.3. Any 4-form T in Λ^4_+ with $dim_{\mathbb{R}}Z_T=6$ determines uniquely an SU(4)-structure. That is, there exists a compatible and positive almost Hermitian structure J on V such that T belongs to λ^4 . The isotropy algebra of T is isomorphic to $\mathfrak{su}(4)$.

Proof. Let $L = Z_T^{\perp}$ be the orthogonal complement of Z_T in \mathcal{S}^+ . Since this is 2-dimensional we get a positive almost Hermitian structure J on V. Now μ_T is completely determined by its restriction to L which gives an element in $S_0^2(L)$ and the fact that T belongs to λ^4 follows from Proposition 6.2. The claim concerning the isotropy algebra follows from Proposition 5.1, (iii) by making use of the above mentioned special isomorphism $\mathfrak{su}(4) \cong \mathfrak{so}(6)$.

This shows how to construct examples of self-dual 4-forms T such that Z_T is of dimension 6. Similarly, from the classification of self-dual 4-forms on \mathbb{R}^8 obtained in [5] one can easily give a geometric description of the cases when Z_T has smaller dimension, but for considerations of time and space we shall not present those here.

6.2. The holonomy algebra. As a convenient intermediary object, we shall make use of the Lie sub-algebra $\mathfrak{g}_T^{*,2}$ of $\mathfrak{g}_T^{*,0} \subseteq A^0$ generated by the sub-set

$$\{[X \,\lrcorner\, T, Y \,\lrcorner\, T] : X,Y \in V\}$$

of A^0 . We point out that a priori $\mathfrak{g}_T^{*,0} \neq \mathfrak{g}_T^{*,2}$.

Lemma 6.2. We have

$$\mathfrak{g}_T^{*,2} = (3+\nu)F \oplus (1-\nu)\iota_T^0$$

Proof. Follows by a straightforward computation based on the fact that $\mathfrak{g}_T^{*,2}$ is generated by the set given in (5.3) and on the equations defining the spaces E and F.

For notational convenience let us set $Q_T^1 = \{X \perp T : X \in V\}$ and also $Q_T^2 = \{X \perp (T\alpha_{12}) : X \in V\}$. Given that $T\alpha_{12} + \alpha_{12}T = 0$ is easily seen that $T\alpha_{12}$ belongs to $\Lambda_+^4(V)$ and hence Q_T^k , k=1,2 are both contained in $\Lambda^3(V)$.

Lemma 6.3. The following hold:

- (i) $[(1-\nu)\iota_T^0, Q_T^1] = Q_T^1$
- (ii) $[(3+\nu)F_{12}, Q_T^1] \subseteq Q_T^1 \oplus Q_T^2$ (iii) $[Q_T^1, Q_T^2] = \mathfrak{g}^{*,2}$

Proof. (i) If α belongs to ι_T^0 and X is in V, an easy computation using essentially that $T\alpha = \alpha T = 0$ and the self-duality of T yields

$$[(1-\nu)\alpha, X \rfloor T] = -2[X \rfloor \alpha, T].$$

(ii) Recall that $T\alpha_{12} + \alpha_{12}T = 0$ and again using the self-duality of T we obtain after a short computation

$$[(3+\nu)\alpha_{12}, X \rfloor T] = 3[X, T\alpha_{12}] - 2[X \rfloor \alpha_{12}, T]$$

for all X in V.

(iii) Because we also have $T\alpha_{12}T = -\frac{1}{2}(1+\nu)\alpha_{12}$ it follows that $T^2\alpha_{12} = \alpha_{12}T^2 = \frac{1}{2}(1+\nu)\alpha_{12}$. Therefore, by using mainly the stability relations and that T is self-dual, we arrive after computing at some length at

$$-4[X \rfloor T, Y \rfloor (T\alpha_{12})] = (TXYT)\alpha_{12} + \alpha_{12}(TYXT) + \frac{1}{2}(1-\nu)(X\alpha_{12}Y + Y\alpha_{12}X)$$

whenever X, Y belong to V. Now using (6.1) it is easily seen that $[T\alpha T, \alpha_{12}] = 0$ for all α in $\Lambda^2(V)$ hence our commutator becomes

$$-4[X \, \lrcorner \, T, Y \, \lrcorner \, (T\alpha_{12})] = -\langle X, Y \rangle (1+\nu)\alpha_{12} + \frac{1}{2}(1-\nu)(X\alpha_{12}Y + Y\alpha_{12}X)$$

for all X, Y in V. On the other hand, given that α_{12} induces a compatible almost complex structure J on V such that $\alpha_{12} = \langle J \cdot, \cdot \rangle$ we actually have

$$X\alpha_{12}Y + Y\alpha_{12}X = 2(X \wedge (Y \perp \alpha_{12}) + Y \wedge (X \perp \alpha_{12})) - 2\langle X, Y \rangle \alpha_{12}$$
$$= 2(X \wedge JY + Y \wedge JX)_0 - \langle X, Y \rangle \alpha_{12}$$

where the subscript indicates orthogonal projection onto Λ_0^2 . Henceforth, our commutator reads finally

$$-4[X \, \lrcorner \, T, Y \, \lrcorner \, (T\alpha_{12})] = -\frac{1}{2} \langle X, Y \rangle (3 + \nu)\alpha_{12} + (1 - \nu)(X \wedge JY + Y \wedge JX)_0$$

for all X,Y in V. Obviously, $(X \wedge JY + Y \wedge JX)_0$ belongs to $\lambda_0^{1,1} = \iota_T^0$ hence $[Q_T^1,Q_T^2] \subseteq \mathfrak{g}_T^{*,2}$ and the equality follows at once when using the linear isomorphism $S_0^2 \to \lambda_0^{1,1}, S \to SJ$.

Theorem 6.1. Let T belong to $\Lambda_+^4(V)$ satisfy $\dim_{\mathbb{R}} Z_T = 6$. Then \mathfrak{g}_T^* is isomorphic to $\mathfrak{so}(6,2)$.

Proof. It now easy to infer from the above that $\mathfrak{g}_T^* = \mathfrak{g}_T^{*,2} \oplus Q_T^1 \oplus Q_T^2$, therefore the result follows.

The proof of Theorem 1.3 is now complete. We end this section by pointing out that in the case above the Clifford multiplication map $\mu: \mathfrak{g}_T^* \to Hom(Z_T^{\perp}, \mathcal{F}^-)$ is no longer surjective.

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APPENDIX A. PROOF OF LEMMA 5.7

Since after use of Lemma 2.3

$$16|T^{2k+1}|^2 = \sum_{q=1}^{p} m_q \lambda_q^{2(2k+1)}$$

for all k in \mathbb{N} we obtain further

(A.1)
$$\frac{1}{4} \sum_{q=1}^{p} m_q \lambda_q^{2(2k+1)} = \frac{1}{2} (\lambda_i^{2(2k+1)} + \lambda_j^{2(2k+1)}) + (\lambda_i \lambda_j)^{2k} (4|T|^2 - \frac{1}{2} (\lambda_i^2 + \lambda_j^2))$$

whenever k belongs to \mathbb{N} .

Case I: $|\lambda_i| \neq |\lambda_i|$

To fix ideas let us assume that $|\lambda_i| < |\lambda_j|$. It follows that

$$\sum_{q=1}^{p} m_q \lim_{k \to \infty} \left(\frac{|\lambda_q|}{|\lambda_j|}\right)^{2(2k+1)} = 2$$

hence $|\lambda_q| \leq |\lambda_j|$ for all $1 \leq q \leq p$ and moreover

$$(A.2) \qquad \sum_{|\lambda_q|=|\lambda_j|} m_q = 2.$$

Therefore

(A.3)
$$\frac{1}{4} \sum_{|\lambda_q| \neq |\lambda_i|}^p m_q \lambda_q^{2(2k+1)} = \frac{1}{2} \lambda_i^{2(2k+1)} + (\lambda_i \lambda_j)^{2k} (4|T|^2 - \frac{1}{2} (\lambda_i^2 + \lambda_j^2))$$

whenever k belongs to N. Further on, after dividing by $(\lambda_i \lambda_j)^{2k+1}$ and taking the limit when $k \to \infty$ we get

$$\frac{1}{4} \sum_{\substack{|\lambda_q| \neq |\lambda_i| \\ |\lambda_q| \neq |\lambda_i|}}^{p} m_q \lim_{k \to \infty} \left(\frac{\lambda_q^2}{|\lambda_i| |\lambda_j|} \right)^{2k+1} = \frac{1}{|\lambda_i| |\lambda_j|} (4|T|^2 - \frac{1}{2} (\lambda_i^2 + \lambda_j^2))$$

leading to $\lambda_q^2 \leq |\lambda_i| |\lambda_j|$ for all $|\lambda_q| \neq |\lambda_j|$ and

(A.4)
$$\sum_{\substack{\lambda_q^2 = |\lambda_i| |\lambda_j| \\ |\lambda_q| \neq |\lambda_j|}} m_q = \frac{4}{|\lambda_i| |\lambda_j|} (4|T|^2 - \frac{1}{2} (\lambda_i^2 + \lambda_j^2)).$$

Therefore, when actualising (A.3) by (A.4) we get

(A.5)
$$\sum_{\substack{\lambda_q^2 \neq |\lambda_i| |\lambda_j| \\ |\lambda_q| \neq |\lambda_j|}}^p m_q \lambda_q^{2(2k+1)} = 2\lambda_i^{2(2k+1)}$$

for all k in \mathbb{N} . We now divide by $\lambda_i^{2(2k+1)}$ and take the limit when $k \to \infty$ to find that

$$\sum_{\substack{|\lambda_q|\neq |\lambda_j|\\\lambda_q^2\neq |\lambda_i\lambda_j|}}^p m_q |\frac{\lambda_q}{\lambda_i}|^{2(2k+1)}=2$$

which implies that $|\lambda_q| \leq |\lambda_i|$ provided that $|\lambda_q| \neq |\lambda_j|$ and $\lambda_q^2 \neq |\lambda_i \lambda_j|$, and also that

(A.6)
$$\sum_{\substack{|\lambda_q| \neq |\lambda_j| \\ \lambda_q^2 \neq |\lambda_i \lambda_j| \\ |\lambda_q| = |\lambda_i|}}^p m_q = 2.$$

When actualising (A.5) by the equation above it follows that the set

$$\{\lambda_q \in \sigma_T : |\lambda_q| \neq |\lambda_j|, |\lambda_q| \neq |\lambda_j|, \lambda_q^2 \neq |\lambda_i\lambda_j|\}$$

is actually empty, in other words $|\lambda_q|$ belongs to $\{|\lambda_i|, |\lambda_j|, \sqrt{|\lambda_i\lambda_j|}\}$ whenever λ_q belongs to σ_T . Moreover, from (A.2) and (A.6) combined with $|\lambda_i| \neq |\lambda_j|$ we get that $m_j \leq 2$ together with $m_i \leq 2$. Also, σ_T always contains eigenvalues λ_q with $\lambda_q^2 = |\lambda_i\lambda_j|$; otherwise, σ_T would be included in $\{\pm\lambda_i, \pm\lambda_j\}$ and moreover (A.5) would imply that either $m_i = 2$ or $m_1 = 1$ and σ_T contains $-\lambda_i$ with multiplicity 1, both of which cannot hold on an 8-dimensional space. Now, by a case by case discussion we shall consider all possibilities.

(i) $m_i = m_j = 2$.

From (A.2) and (A.6) it follows that $-\lambda_i, -\lambda_j$ are not present in σ_T hence the presence of eigenvalues λ_q with $\lambda_q^2 = |\lambda_i| |\lambda_j|$, of which we can have at most 2, leads to the cases $\sigma_T = \{\lambda_i, \lambda_j, \pm \sqrt{|\lambda_i| |\lambda_j|}\}$ or $\sigma_T = \{\lambda_i, \lambda_j, \sqrt{|\lambda_i| |\lambda_j|}, -\sqrt{|\lambda_i| |\lambda_j|}\}$. In all cases, after counting possible multiplicities for the new eigenvalues and using that μ_T is traceless we arrive at $|\lambda_i| = |\lambda_j|$ a contradiction.

(ii) $m_i = 1, m_j = 2$.

In this case $-\lambda_i$, λ_i belong to σ_T each of which with multiplicity 1. Because we have eigenvalues λ_q in σ_T with $\lambda_q^2 = |\lambda_i \lambda_j|$ we have either $\sigma_T = \{\lambda_j, \lambda_i, -\lambda_i, \pm \sqrt{|\lambda_i \lambda_j|}\}$ or $\sigma_T = \{\lambda_j, \lambda_i, -\lambda_i, -\sqrt{|\lambda_i \lambda_j|}, \sqrt{|\lambda_i \lambda_j|}\}$. The latter case cannot be retained because μ_T is trace free whilst in the first case the last eigenvalue has multiplicity 4, as the ambient space is 8-dimensional.

(iii) $m_i = 2, m_i = 1.$

This is completely similar to the case (ii) since λ_i and λ_j , if no ordering is assumed, play dual roles.

(iv) $m_i = m_j = 1$.

Here $\pm \lambda_i, \pm \lambda_j$ belong to σ_T , all of them being simple. The only ways to complete σ_T are $\sigma_T = \{\lambda_i, -\lambda_i, \lambda_j, -\lambda_j, \pm \sqrt{|\lambda_i \lambda_j|}\}$ or $\sigma_T = \{\lambda_i, -\lambda_i, \lambda_j, -\lambda_j, -\sqrt{|\lambda_i \lambda_j|}, \sqrt{|\lambda_i \lambda_j|}\}$ but the first is quickly discarded after using that μ_T is trace free. Concerning the second one, again by the vanishing of the trace and $|\lambda_i| \neq |\lambda_j|$ we find that the multiplicities of the last two eigenvalues can only equal 2.

Case II: $|\lambda_i| = |\lambda_j|$

In this situation it is easily seen that our equation becomes $\sum_{q \neq i,j} m_q \lambda_q^{2(2k+1)} = \lambda_i^{4k} \sum_{q \neq i,j} m_q \lambda_q^2$ for all k in \mathbb{N} . Dividing by λ_i^{4k} and taking the limit when $k \to \infty$ we obtain as above that $|\lambda_q| \leq |\lambda_i|$ for all $q \neq i, j$ and also that $\sum_{\substack{|\lambda_q| = |\lambda_i| \\ q \neq i, j}} m_q \lambda_q^2 = \sum_{q \neq i, j} m_q \lambda_q^2$.

But this implies that $|\lambda_q| = |\lambda_i|$ for all $q \neq i, j$ which is an impossibility for it would imply $\lambda_q = \pm \lambda_i$ hence q = i, j as $\lambda_i + \lambda_j = 0$. Therefore $\sigma_T = \{\lambda_i, \lambda_j\}$ and the claim follows.

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