

Equations for the overlaps of a SIC

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Abstract

We give a holomorphic quartic polynomial in the overlap variables whose zeros on the torus are precisely the Weyl-Heisenberg SICs (symmetric informationally complete positive operator valued measures). By way of comparison, all the other known systems of equations that determine a Weyl-Heisenberg SIC involve variables and their complex conjugates. We also give a related interesting result about the powers of the projective Fourier transform of the group $G = \mathbb{Z}_d \times \mathbb{Z}_d$.

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1 Introduction

Throughout fix the integer $d \geq 2$, and let ω be the primitive d -th root of unity $\omega = e^{\frac{2\pi}{d}i}$. We think of vectors in \mathbb{C}^d as periodic signals on the group \mathbb{Z}_d , and hence index vectors and matrices by elements of \mathbb{Z}_d . A set of d^2 unit vectors (v_j) in \mathbb{C}^d (or the lines that they determine) is said to be **equiangular** if

$$|\langle v_j, v_k \rangle|^2 = \frac{1}{d+1}, \quad j \neq k. \quad (1.1)$$

In quantum information theory, the corresponding rank one orthogonal projections $(v_j v_j^*)$ are said to be a **symmetric informationally complete positive operator valued measure**, or a **SIC** for short. The existence of a SIC for every dimension d is known as *Zauner's conjecture* (from his 1999 thesis, see [Zau10]), or as the *SIC problem*.

There are high precision numerical constructions of SICs [RBKSC04], [SG10], [Sco17], and exact SICs in various dimensions [ACFW18], [GS17]. In all of these constructions, the SIC is a **Weyl-Heisenberg SIC**, i.e., is the orbit $(\rho(g)v)_{g \in G}$ of a *fiducial vector* v under the unitary irreducible projective representation $\rho : G \rightarrow \mathcal{U}(\mathbb{C}^{\mathbb{Z}_d})$ of $G = \mathbb{Z}_d \times \mathbb{Z}_d$ with Schur multiplier α given by

$$\rho_{jk} = \rho((j, k)) = S^j \Omega^k, \quad \alpha((j_1, j_2), (k_1, k_2)) = \omega^{j_2 k_1}, \quad (1.2)$$

where S is the cyclic shift matrix $S_{jk} := \delta_{j, k+1}$ and Ω is the diagonal (modulation) matrix $\Omega_{jk} := \omega^j \delta_{jk}$. In this case, the equiangularity condition (1.1) becomes

$$|\langle S^j \Omega^k v, v \rangle|^2 = \frac{1}{d+1}, \quad (j, k) \neq (0, 0). \quad (1.3)$$

In this paper, we consider equations in the variables

$$c_{jk} = \langle S^j \Omega^k v, v \rangle = \text{trace}(v v^* S^j \Omega^k), \quad (j, k) \in \mathbb{Z}_d \times \mathbb{Z}_d, \quad (1.4)$$

which determine a (Weyl-Heisenberg) SIC. These variables (or scalar multiples of them) are called the **overlaps** of the SIC. They depend only on the fiducial projector $P = v v^*$. The original attempts to find numerical and exact SIC fiducials (using Groebner basis methods) involved polynomial equations in the variables v_0, \dots, v_{d-1} and $\bar{v}_0, \dots, \bar{v}_{d-1}$, such as the equiangularity condition (1.3), the equations (see [BW07], [Kha08], [ADF14])

$$\sum_{r \in \mathbb{Z}_d} v_r \bar{v}_{r+s} \bar{v}_{r+t} v_{r+s+t} = \begin{cases} 0, & s, t \neq 0; \\ \frac{1}{d+1}, & s \neq 0, t = 0, \quad s = 0, t \neq 0; \\ \frac{2}{d+1}, & (s, t) = (0, 0), \end{cases} \quad (1.5)$$

and the variational characterisation (used for finding numerical SICs)

$$\frac{1}{d^2} \sum_{(j,k) \in \mathbb{Z}_d^2} |\langle S^j \Omega^k v, v \rangle|^4 = \frac{2}{d(d+1)} \|v\|^4, \quad \|v\|^2 = 1. \quad (1.6)$$

More recent exact constructions of SICs [ACFW18] have been in the overlap variables c_{jk} (utilising a natural Galois action on them). Clearly the c_{jk} giving a SIC fiducial projector vv^* via (1.4) must satisfy

$$c_{00} = \|v\|^2 = 1, \quad |c_{jk}|^2 = |\langle S^j \Omega^k v, v \rangle|^2 = \frac{1}{d+1}, \quad (j, k) \neq (0, 0), \quad (1.7)$$

and also, by the rule $\Omega^k S^j = \omega^{jk} S^j \Omega^k$,

$$c_{jk} = \overline{\langle v, S^j \Omega^k v \rangle} = \overline{\langle \Omega^{-k} S^{-j} v, v \rangle} = \overline{\langle \omega^{jk} S^{-j} \Omega^{-k} v, v \rangle} = \omega^{-jk} \overline{c_{-j, -k}}. \quad (1.8)$$

These conditions on the overlap variables c_{jk} are not enough to guarantee that they come from a fiducial projector vv^* (and hence prove Zauner's conjecture).

In Section 2, we define a linear operator T , which is an example of the projective Fourier transform, which allows us to reconstruct the fiducial projector as $vv^* = Tc$ from a suitable $c = (c_{jk})$. We prove that in addition to (1.7) and (1.8), the simple condition

$$\text{trace}((Tc)^4) = 1$$

ensures that a c gives a SIC fiducial (Theorem 2.1). We then give some examples, and describe the action of the Clifford group on the SIC fiducials give by overlaps c .

In Section 3, we give some interesting properties of T , i.e., the projective Fourier transform of $G = \mathbb{Z}_d \times \mathbb{Z}_d$. In particular, we show that $(\sqrt{dT})^{6d} = (-1)^{\frac{1}{2}d(d-1)} I$, and a variant has order $4d$. To our knowledge, this is only the second example of a Fourier transform of finite order, after the (discrete) Fourier transform for a finite abelian group $G = \mathbb{Z}_d$ (which satisfies $F^4 = I$).

In Section 4, we give another system of equations in the overlaps c that determine a SIC. These involve the symbol (z -transform) of the rows of c . The symbols for c giving a SIC turn out to have interesting Riesz-type factorisation properties. We use these to describe the (sporadic) SICs for $d = 3$, which are parametrised by a hypocycloid.

2 The reconstruction operator

Since the ρ of (1.2) is a unitary irreducible projective representation of dimension d , it follows that $(\rho(g))_{g \in G}$ is a tight frame (called a nice error frame with index group G [CW17]) for the $d \times d$ matrices with the Frobenius inner product

$$\langle A, B \rangle := \text{trace}(AB^*) = \sum_{j,k} a_{jk} \overline{b_{jk}},$$

i.e.,

$$A = \frac{d}{|G|} \sum_{g \in G} \langle A, \rho(g) \rangle \rho(g), \quad \forall A \in \mathbb{C}^{d \times d}. \quad (2.9)$$

In this particular case, $(\rho(g))_{g \in G} = (S^j \Omega^k)$ is an orthogonal basis. Taking $A = vv^*$ above gives the following formula for reconstruction from the overlaps $c_{jk} = \langle S^j \Omega^k v, v \rangle$

$$vv^* = \frac{d}{d^2} \sum_{j,k} \langle vv^*, S^{-j} \Omega^{-k} \rangle S^{-j} \Omega^{-k} = \frac{1}{d} \sum_{j,k} \omega^{jk} c_{jk} S^{-j} \Omega^{-k} = \frac{1}{d} \sum_{j,k} c_{jk} (S^j \Omega^k)^*,$$

since $\omega^{jk}S^{-j}\Omega^{-k} = (S^j\Omega^k)^*$, and $(S^{-j}\Omega^{-k})^* = \omega^{jk}S^j\Omega^k$ gives

$$\langle vv^*, S^{-j}\Omega^{-k} \rangle = \text{trace}(vv^*(S^{-j}\Omega^{-k})^*) = \text{trace}(v^*\omega^{jk}S^j\Omega^k v) = \omega^{jk}\langle S^j\Omega^k v, v \rangle = \omega^{jk}c_{jk}.$$

Motivated by this, we define a linear map $T : \mathbb{C}^{\mathbb{Z}_d \times \mathbb{Z}_d} \rightarrow \mathbb{C}^{\mathbb{Z}_d \times \mathbb{Z}_d}$ by

$$Tc := \frac{1}{d} \sum_{j,k} c_{jk} (S^j\Omega^k)^* = \frac{1}{d} \sum_{j,k} \omega^{jk} c_{jk} S^{-j}\Omega^{-k}. \quad (2.10)$$

This can be viewed as the α -Fourier transform of [Wal20] (for a Schur multiplier α) which is a map $F_\alpha : \mathbb{C}^G \rightarrow \bigoplus_\rho \mathbb{C}^{d_\rho \times d_\rho}$, where ρ counts over the irreducible projective representations of G with multiplier α (and dimension d_ρ). Here $G = \mathbb{Z}_d \times \mathbb{Z}_d$ has just one such representation, the ρ of (1.2), and F_α of $\nu = c \in \mathbb{C}^G = \mathbb{C}^{\mathbb{Z}_d \times \mathbb{Z}_d}$ at the unitary representation ρ is

$$(F_\alpha \nu)_\rho = \sum_{g \in G} \nu(g) \rho(g)^* = \sum_{j,k} c_{jk} (S^j\Omega^k)^* = d(Tc).$$

Thus T is the projective Fourier transform for the group $G = \mathbb{Z}_d \times \mathbb{Z}_d$. For this particular group, we can view the image of a vector in \mathbb{C}^G as being in $\mathbb{C}^G = \mathbb{C}^{\mathbb{Z}_d \times \mathbb{Z}_d}$, and as a result it is natural to consider powers of the Fourier transform. The only other case that we know of where this can be done is for the ordinary representations of a finite abelian group (where the representations give the character group \hat{G} , which can be identified with G). In this case the (discrete) Fourier transform has order 4.

We now use T to characterise those vectors (matrices) $c \in \mathbb{C}^G = \mathbb{C}^{\mathbb{Z}_d \times \mathbb{Z}_d}$ which give a fiducial projector $vv^* = Tc$.

Lemma 2.1 *Let T be given by (2.10). Suppose that $c = (c_{jk}) \in \mathbb{C}^{\mathbb{Z}_d \times \mathbb{Z}_d}$ satisfies*

- (i) $c_{jk} = \omega^{-jk} \overline{c_{-j,-k}}$
- (ii) $c_{00} = 1$
- (iii) $|c_{jk}|^2 = \frac{1}{d+1}$, $(j,k) \neq (0,0)$.

Then Tc is Hermitian, and its eigenvalues $\lambda_1, \dots, \lambda_d$ satisfy

$$\sum_j \lambda_j = 1, \quad \sum_j \lambda_j^2 = 1, \quad \sum_{j \neq k} \lambda_j \lambda_k = 0,$$

i.e., its characteristic polynomial has the form

$$p_{Tc}(\lambda) = \lambda^d - \lambda^{d-1} + 0\lambda^{d-2} + a_{d-3}\lambda^{d-3} + \dots + a_1\lambda + a_0.$$

Proof: Firstly, observe (i) implies that Tc is Hermitian, since

$$(Tc)^* = \frac{1}{d} \sum_{j,k} \omega^{jk} c_{-j,-k} S^j\Omega^k = \frac{1}{d} \sum_{j,k} \omega^{(-j)(-k)} c_{j,k} S^{-j}\Omega^{-k} = Tc.$$

Since $\text{trace}(S^j \Omega^k) = 0$, $(j, k) \neq (0, 0)$, we calculate using (ii) that

$$\sum_j \lambda_j = \text{trace}(Tc) = \frac{1}{d} \sum_{j,k} \omega^{jk} c_{j,k} \text{trace}(S^{-j} \Omega^{-k}) = \frac{1}{d} \omega^0 c_{00} \text{trace}(I) = c_{00} = 1.$$

The so called 2-trace $\sum_{j \neq k} \lambda_j \lambda_k$ of Tc is equal to $\{(\text{trace}(Tc))^2 - \text{trace}((Tc)^2)\}/2$. Since Tc is Hermitian, $\text{trace}((Tc)^2) = \langle Tc, (Tc)^* \rangle = \langle Tc, Tc \rangle$, and by the orthogonality of the $\rho_{jk} = S^j \Omega^k$, we calculate

$$\langle Tc, Tc \rangle = \frac{1}{d^2} \sum_{j,k} |c_{j,k}|^2 \langle \rho_{jk}^*, \rho_{jk}^* \rangle = \frac{1}{d} \sum_{j,k} |c_{j,k}|^2.$$

Now by (ii) and (iii),

$$\text{trace}((Tc)^2) = \langle Tc, Tc \rangle = \frac{1}{d} \sum_{j,k} |c_{j,k}|^2 = \frac{1}{d} \left(1 + (d^2 - 1) \frac{1}{d+1} \right) = 1,$$

and so $\sum_{j \neq k} \lambda_j \lambda_k = \{(\text{trace}(Tc))^2 - \text{trace}((Tc)^2)\}/2 = (1 - 1)/2 = 0$. □

Theorem 2.1 *Let T be given by (2.10). Then a matrix $c = (c_{jk}) \in \mathbb{C}^{\mathbb{Z}_d \times \mathbb{Z}_d}$ determines a fiducial projector for a Weyl-Heisenberg SIC by $vv^* = Tc$ if and only if*

(i) $c_{jk} = \omega^{-jk} \overline{c_{-j, -k}}$

(ii) $c_{00} = 1$

(iii) $|c_{jk}|^2 = \frac{1}{d+1}$, $(j, k) \neq (0, 0)$

(iv) $\text{trace}((Tc)^4) = 1$

Moreover this fiducial satisfies $\langle S^j \Omega^k v, v \rangle = c_{jk}$, and if $v_0 \neq 0$, then v is given by

$$v = \frac{1}{d v_0} \bar{c} \begin{pmatrix} 1 \\ 1 \\ \vdots \\ 1 \end{pmatrix}.$$

Proof: By Lemma 2.1, the eigenvalues of the Hermitian matrix Tc satisfy

$$\sum_j \lambda_j = 1, \quad \sum_j \lambda_j^2 = 1, \quad \sum_{j \neq k} \lambda_j \lambda_k = 0,$$

so that $0 \leq \lambda_j^2 \leq 1$. Thus $\lambda_j^4 \leq \lambda_j^2$, with equality if and only if $\lambda_j^2 \in \{0, 1\}$. But

$$\sum_j \lambda_j^4 = \text{trace}((Tc)^4) = 1 = \text{trace}((Tc)^2) = \sum_j \lambda_j^2,$$

so that $\lambda_j^2 \in \{0, 1\}$, $\forall j$, and we must have $\lambda_j = 1$ for some j , and $\lambda_j = 0$ for all others, i.e., Tc is rank one, say

$$Tc = \frac{1}{d} \sum_{j,k} c_{jk} \rho_{jk}^* = vv^*, \quad v \in \mathbb{C}^{\mathbb{Z}_d}.$$

Since $\{\rho_{jk}\}$ is orthogonal, taking the inner product of the above with $\rho_{jk} = S^j \Omega^k$ gives

$$c_{jk} = \frac{1}{d} c_{jk} \langle \rho_{jk}^*, \rho_{jk}^* \rangle = \langle vv^*, \rho_{jk}^* \rangle = \text{trace}(vv^* S^j \Omega^k) = \text{trace}(v^* S^j \Omega^k v) = \langle S^j \Omega^k v, v \rangle.$$

Finally, with e_j the standard basis vectors, we calculate

$$\begin{aligned} j\text{-th entry of } \bar{c} \begin{pmatrix} 1 \\ 1 \\ \vdots \\ 1 \end{pmatrix} &= \sum_k \bar{c}_{jk} = \sum_k \langle v, S^j \Omega^k v \rangle = \langle v, S^j \left(\sum_k \Omega^k \right) v \rangle \\ &= \langle v, S^j \begin{pmatrix} d & 0 & \cdots & 0 \\ 0 & 0 & \cdots & 0 \\ \vdots & \vdots & & \vdots \\ 0 & 0 & \cdots & 0 \end{pmatrix} v \rangle = \langle v, S^j d v_0 e_0 \rangle = d \bar{v}_0 \langle v, e_j \rangle = d \bar{v}_0 v_j. \end{aligned}$$

□

From the proof, we see that (iv) can be replaced by various equivalent conditions, e.g.,

$$(iv)' \text{ The characteristic polynomial of } Tc \text{ has the form } P_{Tc}(\lambda) = \lambda^d - \lambda^{d-1}$$

$$(iv)'' \text{ trace}((Tc)^j) = 1, \quad j = 1, 2, \dots$$

since given (i), (ii), (ii),

$$(iv)'' \implies (iv) \implies Tc \text{ has eigenvalues } 1, 0, \dots, 0 \iff (iv)' \implies (iv)''.$$

By condition (ii), we may set $c_{00} = 1$, to obtain the following characterisation.

Corollary 2.1 *The overlaps of a Weyl-Heisenberg SIC are precisely the zeros of the polynomial $\text{trace}((Tc)^4) = 1$ on the torus*

$$|c_{jk}| = \frac{1}{\sqrt{d+1}}, \quad (j, k) \neq (0, 0).$$

The condition (i) allows further variables c_{jk} to be eliminated. When d is odd, half of the $(d^2 - 1)$ variables c_{jk} , $(j, k) \neq (0, 0)$, can be eliminated. For d even, half of the $d^2 - 4$ variables c_{jk} , $(j, k) \notin \{0, \frac{d}{2}\}^2$, can be eliminated, and

$$c_{\frac{d}{2}, 0} = \overline{c_{\frac{d}{2}, 0}}, \quad c_{0, \frac{d}{2}} = \overline{c_{0, \frac{d}{2}}}, \quad c_{\frac{d}{2}, \frac{d}{2}} = (-1)^{\frac{d}{2}} \overline{c_{\frac{d}{2}, \frac{d}{2}}}, \quad (2.11)$$

so that $c_{0, \frac{d}{2}}, c_{\frac{d}{2}, 0} \in \mathbb{R}$, and $c_{\frac{d}{2}, \frac{d}{2}}$ is in \mathbb{R} for $\frac{d}{2}$ even, and is in $i\mathbb{R}$ for $\frac{d}{2}$ odd.

Example 2.1 For $d = 2$, the conditions (2.11) of (i) give $c_{01} \in \mathbb{R}$, $c_{10} \in \mathbb{R}$, $c_{11} \in i\mathbb{R}$. Hence imposing the conditions (ii) and (iii), we have eight possibilities

$$c_{00} = 1, \quad c_{01} = \pm \frac{1}{\sqrt{3}}, \quad c_{10} = \pm \frac{1}{\sqrt{3}}, \quad c_{11} = \pm i \frac{1}{\sqrt{3}}. \quad (2.12)$$

Taking the ‘+’ choice above gives

$$Tc = T \begin{pmatrix} 1 & \frac{1}{\sqrt{3}} \\ \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{3}}i \end{pmatrix} = \frac{1}{2} \left(I + \frac{1}{\sqrt{3}}(S + \Omega) + \frac{i}{\sqrt{3}}S\Omega \right) = \frac{1}{2\sqrt{3}} \begin{pmatrix} \sqrt{3} + 1 & 1 - i \\ 1 + i & \sqrt{3} - 1 \end{pmatrix},$$

which satisfies $\text{trace}((Tc)^4) = \text{trace}(Tc) = 1$, and so gives a Weyl-Heisenberg SIC

$$v = \frac{1}{\sqrt{2}\sqrt{\sqrt{3} + 1}} \begin{pmatrix} \sqrt{3} + 1 \\ 1 + i \end{pmatrix}.$$

In fact all eight choices give SICs which are equivalent, as we now explain.

The group generated by S and Ω is called the **Heisenberg group**, and its normaliser in the unitary matrices is the **Clifford group**. Indeed, if a $a \in C(d)$, then

$$a\rho_{jk}a^{-1} = z_a(j, k)\rho_{\psi_a(j, k)}, \quad \forall (j, k) \in \mathbb{Z}_d^2,$$

where ψ_a is matrix multiplication by an element of $SL_2(\mathbb{Z}_d)$. The Clifford group $C(d)$ maps SIC fiducials to SIC fiducials, via the action

$$a \cdot (vv^*) := (av)(av)^* = a(vv^*)a^{-1}, \quad a \in C(d).$$

The induced action on the overlaps of the fiducial is given by

$$\begin{aligned} (a \cdot c)_{jk} &= \text{trace}(a(vv^*)a^{-1}S^j\Omega^k) = \langle a^{-1}S^j\Omega^kav, v \rangle = \langle z_{a^{-1}}(j, k)\rho_{\psi_{a^{-1}}(j, k)}v, v \rangle \\ &= z_{a^{-1}}(j, k)c_{\psi_{a^{-1}}(j, k)}. \end{aligned}$$

In [BW19], it is shown that the Clifford group is generated by the scalar matrices, S , Ω , the Fourier transform F and the Zauner matrix Z , where

$$F_{jk} := \frac{1}{\sqrt{d}}\omega^{jk}, \quad Z_{jk} := \zeta^{d-1}\mu^{j(j+d)}, \quad \mu := e^{\frac{2\pi i}{2d}}, \quad \zeta := e^{\frac{2\pi i}{24}}.$$

For these (see [Wal18])

$$(S^a\Omega^b \cdot c)_{jk} = \omega^{ak-bj}c_{jk}, \quad (F \cdot c)_{jk} = \omega^{-jk}c_{k,-j}, \quad (Z \cdot c)_{jk} = \mu^{j(j+d-2k)}c_{k-j,-j}.$$

When a (Weyl-Heisenberg) SIC fiducial vv^* is known, there is always appears to be one which is given by an eigenvector v of Z (indeed these are often searched for directly). Correspondingly, the overlaps satisfy $Z \cdot c = c$, i.e., the equations

$$\mu^{j(j+d-2k)}c_{k-j,-j} = c_{jk},$$

which allows a further reduction of the variables c_{jk} .

3 Properties of the projective Fourier transform

Here we consider some properties of $T : \mathbb{C}^{\mathbb{Z}_d \times \mathbb{Z}_d} \rightarrow \mathbb{C}^{\mathbb{Z}_d \times \mathbb{Z}_d}$ given by (2.10), i.e., the projective Fourier transform of $G = \mathbb{Z}_d \times \mathbb{Z}_d$. It follows from the Plancherel formula for projective representations [Wal20], or (2.9) that $\sqrt{d}T$ is unitary. Indeed, (2.9) can be written as $I = T\Lambda$, where $\Lambda : \mathbb{C}^G \rightarrow \mathbb{C}^G : A \mapsto (\langle A, \rho(g) \rangle)_{g \in G}$ satisfies

$$\langle A, A \rangle = \frac{1}{d} \sum_{g \in G} |\langle A, \rho(g) \rangle|^2 = \frac{1}{d} \langle \Lambda A, \Lambda A \rangle,$$

so that $\frac{1}{\sqrt{d}}\Lambda$ is unitary, and hence $\sqrt{d}T$ is unitary.

We now show that $\sqrt{d}T$ has finite order ($6d$ or $12d$), i.e., the projective Fourier transform for ρ of (1.2) has finite order (Theorem 3.1). To do this, we need a technical lemma (Lemma 3.1), based on the Zauner matrix Z (of order 3), which can be factored

$$Z = \zeta^{d-1}RF, \quad \zeta := e^{\frac{2\pi i}{24}}, \quad (R)_{jk} = \mu^{j(j+d)}\delta_{jk}, \quad \mu := e^{\frac{2\pi i}{2d}},$$

where F is the Fourier matrix, and R is diagonal. The *strong* form of Zauner's conjecture is that there is a SIC fiducial which is an eigenvector of Z , for every dimension d .

Lemma 3.1 *For any d , we have that*

$$(R^2F)^2 = \zeta^{-6(d-1)}(RF)R^{-2}(RF)^{-1}, \quad \zeta = e^{\frac{2\pi i}{24}},$$

and, in particular

$$(R^2F)^{2d} = (-1)^{\frac{1}{2}d(d-1)}I.$$

Proof: Write $Z = cRF$, $c = \zeta^{d-1}$. Since $Z^3 = I$ and $F^4 = I$, we have

$$\begin{aligned} R^2F &= R(RF)^2(RF)^{-1} = R(\bar{c}Z)^2(RF)^{-1} = \bar{c}^2RZ^{-1}(RF)^{-1} = \bar{c}^3RF^{-1}R^{-1}(RF)^{-1} \\ &= \bar{c}^3(RF)(F^2R^{-1})(RF)^{-1}. \end{aligned}$$

Since the permutation matrix F^2 commutes with R (or any power of R), we have

$$(R^2F)^2 = \bar{c}^6(RF)(F^2R^{-1})^2(RF)^{-1} = \bar{c}^6(RF)R^{-2}(RF)^{-1},$$

where $\bar{c}^6 = \zeta^{-6(d-1)}$. Since $R^{2d} = I$, we obtain

$$(R^2F)^{2d} = \bar{c}^{6d}(RF)R^{-2d}(RF)^{-1} = \bar{c}^{6d}I, \quad \bar{c}^{6d} = \zeta^{-6d(d-1)} = (-1)^{\frac{1}{2}d(d-1)},$$

which completes the proof. □

Theorem 3.1 *The reconstruction operator T of (2.10) has finite order, i.e.,*

$$(\sqrt{d}T)^{6d} = (-1)^{\frac{1}{2}d(d-1)}I.$$

Proof: We consider T with respect to the standard basis $E_{jk} = e_j e_k^*$ for matrices, ordered so that the coordinates of c have the block structure $[c] = (c_0, \dots, c_{d-1})^T$, where c_j is the j -th column of the matrix c (this is the order of matlab's `reshape(c, d, 2, 1)`).

The (j, k) -block A_{jk} of the (block) matrix representation $[\sqrt{d}T]$ of $\sqrt{d}T$ is given by

$$\begin{aligned} A_{jk}v &= j\text{-th column of } \sqrt{d}T([0 \dots 0, v, 0 \dots 0]) \quad (v \text{ is the } k\text{-th column}) \\ &= \frac{1}{\sqrt{d}} \sum_{a,b} [0 \dots 0, v, 0 \dots 0]_{ab} (S^a \Omega^b)^* e_j = \frac{1}{\sqrt{d}} \sum_a v_a \Omega^{-k} S^{-a} e_j \\ &= \frac{1}{\sqrt{d}} \sum_a (\Omega^{-k} P_{-1} S^{-j} e_a) v_a = \frac{1}{\sqrt{d}} (\Omega^{-k} P_{-1} S^{-j}) v, \end{aligned}$$

so that

$$A_{jk} = \frac{1}{\sqrt{d}} \Omega^{-k} P_{-1} S^{-j}, \quad (A_{jk})_{ab} = \frac{1}{\sqrt{d}} \omega^{-ak} \delta_{a, j-b}.$$

The (j, k) -block B_{jk} of $[\sqrt{d}T]^2$ is given by

$$\begin{aligned} (B_{jk})_{ab} &= \left(\sum_r A_{jr} A_{rk} \right)_{ab} = \sum_r \sum_t (A_{jr})_{at} (A_{rk})_{tb} = \frac{1}{d} \sum_r \sum_t \omega^{-ar} \delta_{a, j-t} \omega^{-tk} \delta_{t, r-b} \\ &= \frac{1}{d} \omega^{-a(j-a+b)} \omega^{-(j-a)k} = \frac{1}{d} \omega^{a^2 - aj + ak - jk - ab}. \end{aligned}$$

The (j, k) -block C_{jk} of $[\sqrt{d}T]^3$ is given by

$$\begin{aligned} (C_{jk})_{ab} &= \left(\sum_r B_{jr} A_{rk} \right)_{ab} = \sum_r \sum_t (B_{jr})_{at} (A_{rk})_{tb} = \sum_r \sum_t (B_{jr})_{at} (A_{rk})_{tb} \\ &= \frac{1}{d\sqrt{d}} \sum_r \sum_t \omega^{a^2 - aj + ar - jr - at} \omega^{-tk} \delta_{t, r-b} = \frac{1}{d\sqrt{d}} \sum_r \omega^{a^2 - aj + ar - jr - a(r-b)} \omega^{-(r-b)k} \\ &= \frac{1}{d\sqrt{d}} \omega^{a^2 - aj + ab + bk} \sum_r \omega^{-r(j+k)} = \frac{1}{\sqrt{d}} \omega^{a^2 - aj + ab + bk} \delta_{j, -k} = (R^2 \Omega^{-j} F \Omega^k)_{ab} \delta_{j, -k}, \end{aligned}$$

so that

$$C_{jk} = \begin{cases} 0, & k \neq -j; \\ R^2 \Omega^{-j} F \Omega^{-j}, & k = -j. \end{cases}$$

It therefore follows, that $[\sqrt{d}T]^6$ is block diagonal, with diagonal blocks

$$Q_{jj} = C_{j, -j} C_{-j, j} = (R^2 \Omega^{-j} F \Omega^{-j})(R^2 \Omega^j F \Omega^j) = \Omega^{-j} (R^2 F)^2 \Omega^j.$$

Thus $[\sqrt{d}T]^{6d}$ is block diagonal, and, by Lemma 3.1, its diagonal blocks simplify to

$$\Omega^{-j} (R^2 F)^{2d} \Omega^j = \Omega^{-j} (-1)^{\frac{1}{2}d(d-1)} I \Omega^j = (-1)^{\frac{1}{2}d(d-1)} I,$$

i.e., $[(\sqrt{d}T)^{6d}] = [(-1)^{\frac{1}{2}d(d-1)} I]$. □

Since the projective representation (1.2) of $\mathbb{Z}_d \times \mathbb{Z}_d$ is *not* an ordinary representation, there is no canonical presentation of the projective Fourier transform at ρ , as with the Fourier transform for \mathbb{Z}_d , which gives F (of order 4), by taking $\alpha = 1$. Indeed, one could take $\tilde{\rho}((j, k)) = b_{jk} S^j \Omega^k$, for any unit scalars b_{jk} , with a corresponding $\tilde{\alpha}$ -transform (reconstruction operator)

$$\tilde{T}c := \frac{1}{d} \sum_{j,k} c_{jk} (b_{jk} S^j \Omega^k)^* = \frac{1}{d} \sum_{j,k} c_{jk} \overline{b_{jk}} \omega^{jk} S^{-j} \Omega^{-k}.$$

For a general choice for b_{jk} , $\sqrt{d}\tilde{T}$ is again unitary, but not of finite order. During our investigation, we came across various choices giving operators of finite order, in particular

$$Lc := \frac{1}{d} \sum_{j,k} c_{jk} (\omega^{jk} S^j \Omega^k)^* = \frac{1}{d} \sum_{j,k} c_{jk} S^{-j} \Omega^{-k}. \quad (3.13)$$

It can be shown that L has the compact form

$$Lc = F^*(F \circ (F^* c F^*)), \quad (3.14)$$

where \circ is the Hadamard product. From this, we obtain the following.

Theorem 3.2 *The operator L of (3.13) satisfies*

$$(\sqrt{d}L)^4 c = (F^* R^{-2} F) c \quad (\text{matrix multiplication}),$$

and hence, since $R^{2d} = I$, we have

$$(\sqrt{d}L)^{4d} = I.$$

Proof: We first verify the compact form (3.14),

$$(Lc)_{ab} = \frac{1}{d} \sum_{j,k} c_{jk} (S^{-j} \Omega^{-k})_{ab} = \frac{1}{d} \sum_{j,k} c_{jk} \omega^{-kb} \delta_{a,b-j} = \frac{1}{d} \sum_k c_{b-a,k} \omega^{-kb},$$

$$\begin{aligned} (F^*(F \circ (F^* c F^*)))_{ab} &= \sum_{r,t,k} (F^*)_{ar} (F)_{rb} (F^*)_{rt} c_{tk} (F^*)_{kb} \\ &= \frac{1}{d^2} \sum_{r,t,k} \omega^{-ar+rb-rt} c_{tk} \omega^{-kb} = \frac{1}{d} \sum_{t,k} \delta_{t,b-a} c_{tk} \omega^{-kb} = \frac{1}{d} \sum_k c_{b-a,k} \omega^{-kb}. \end{aligned}$$

Define the operation $\tilde{A} = P_{1-} A P_{-1}$ of conjugation by the permutation matrix $P_{-1} = F^2$ of order 2. This distributes over matrix multiplication, the Hadamard product, leaving F (and its powers) unchanged, so that $Lc = F^*(F \circ (F \tilde{c} F))$, and

$$\begin{aligned} L^2 c &= F^*(F \circ (F[F^*(F \circ (F \tilde{c} F))]F)) = F^*(F \circ (F[F^*(F \circ (F c F))]F)) \\ &= F^*([F \circ ([F \circ (F c F)]F)]F)F^* = F^* M^2 (F c F) F^*, \end{aligned}$$

where

$$\begin{aligned}
Mc &:= (F \circ c)F, & (Mc)_{jk} &= \sum_t (F \circ c)_{jt} F_{tk} = \frac{1}{d} \sum_t \omega^{jt} c_{jt} \omega^{tk} = \frac{1}{d} \sum_t \omega^{(j+k)t} c_{jt}, \\
(M^2c)_{jk} &= \frac{1}{d} \sum_t \omega^{(j+k)t} (Mc)_{jt} = \frac{1}{d} \sum_t \omega^{(j+k)t} \frac{1}{d} \sum_r \omega^{(j+t)r} c_{jr} = \frac{1}{d} \omega^{-j(j+k)} c_{j, -(j+k)}, \\
(M^4c)_{jk} &= \frac{1}{d^2} \omega^{-j(j+k)} \omega^{-j(-k)} c_{jk} = \frac{1}{d^2} \mu^{-2j(j+d)} c_{jk} = \left(\frac{1}{d^2} R^{-2}c \right)_{jk}.
\end{aligned}$$

Thus, $M^4c = \frac{1}{d^2} R^{-2}c$, which gives

$$\begin{aligned}
(\sqrt{d}L)^4c &= d^2 F^* M^2 (F [F^* M^2 (FcF) F^*] F) F^* = d^2 F^* M^4 (FcF) F^* \\
&= F^* (R^{-2} FcF) F^* = F^* R^{-2} Fc,
\end{aligned}$$

and $(\sqrt{d}L)^{4d}c = (F^* R^{-2} F)^d c = F^* R^{-2d} Fc = c$. \square

4 Equivalent equations for Heisenberg frames

In this section, we give another condition that ensures Tc has rank one, which leads to a set of equations for c which express in terms of polynomials $p_j(z)$ which are z -transforms of the rows of c . These polynomials $p_j(z)$ have interesting Riesz-type factorisation properties, which we use to find a solution for $d = 4$.

We use the following condition which ensures that a matrix $A \in \mathbb{C}^{d \times d}$ has rank one.

Lemma 4.1 $A = vv^*$ for some $v \in \mathbb{C}^d$ with $v_m \neq 0$ if and only if $a_{mm} > 0$ and

$$A = \frac{1}{a_{mm}} \begin{bmatrix} a_{0m} \\ a_{1m} \\ a_{2m} \\ \vdots \end{bmatrix} \begin{bmatrix} a_{0m} \\ a_{1m} \\ a_{2m} \\ \vdots \end{bmatrix}^*. \quad (4.15)$$

Proof: First suppose that $A = vv^*$ for such a v . Then

$$\overline{v_m}v = Ae_m = \begin{bmatrix} a_{0m} \\ a_{1m} \\ a_{2m} \\ \vdots \end{bmatrix}, \quad |v_m|^2 = (\overline{v_m}v)_m = a_{mm},$$

so that $a_{mm} > 0$, and (4.15) holds since $(\overline{v_m}v)(\overline{v_m}v)^* = |v_m|^2(vv^*)$.

Conversely, suppose that (4.15) holds with $a_{mm} > 0$, then clearly $A = vv^*$ for

$$v := \frac{1}{\sqrt{a_{mm}}} \begin{bmatrix} a_{0m} \\ a_{1m} \\ a_{2m} \\ \vdots \end{bmatrix}, \quad v_m = \sqrt{a_{mm}}.$$

\square

In particular, $Tc = vv^*$ for some $v \in \mathbb{C}^d$ with $v_m \neq 0$ if and only if $(Tc)_{mm} > 0$ and

$$(Tc)_{jk} = \frac{(Tc)_{jm} \overline{(Tc)_{km}}}{(Tc)_{mm}}. \quad (4.16)$$

We now express (4.16) in terms of the following z -transform.

Definition 4.1 For $j = 0, 1, \dots, d-1$, the j -th **symbol** of c is defined to be the polynomial

$$p_j(z) := \sum_r c_{-j,r} (\omega^j z)^{-r}.$$

This is the z -transform of the j -th row of the matrix $(\omega^{jk} c_{-j,-k})$, since

$$\sum_k \omega^{jk} c_{-j,-k} z^k = \sum_r \omega^{-jr} c_{-j,r} z^{-r} = \sum_r c_{-j,r} (\omega^j z)^{-r}.$$

We think of $p_j(z)$ as being defined only on $z^d = 1$, since each polynomial of degree d is uniquely determined by its values at the d -th roots of unity. Clearly, we can recover c from the d^2 values $p_j(\omega^k)$, $j, k = 0, \dots, d-1$. Using (??), we calculate

$$p_{j-k}(\omega^k) = \sum_r c_{k-j,r} (\omega^{j-k} \omega^k)^{-r} = \sum_r c_{k-j,r} (\omega^j)^{-r} = d(Tc)_{jk}.$$

Hence (4.16) can be expressed as follows.

Theorem 4.1 $Tc = vv^*$ for $v \in \mathbb{C}^d$ with $v_m \neq 0$ if and only if the symbols of c satisfy

$$p_0(\omega^m) > 0, \quad p_{j-k}(\omega^k) = \frac{p_{j-m}(\omega^m) \overline{p_{k-m}(\omega^m)}}{p_0(\omega^m)}, \quad j, k = 0, \dots, d-1. \quad (4.17)$$

The symbols corresponding to a solution have interesting Riesz-type factorisation properties, which, for simplicity, we illustrate when $m = 0$.

Corollary 4.1 $Tc = vv^*$ for $v \in \mathbb{C}^d$ with $v_0 \neq 0$ if and only if the symbols of c satisfy

$$p_0(1) > 0, \quad p_{j-k}(\omega^k) = \frac{p_j(1) \overline{p_k(1)}}{p_0(1)}, \quad j, k = 0, \dots, d-1. \quad (4.18)$$

Moreover, these have the factorisations

$$|p_j(z)|^2 = p_0(z) p_0(\omega^j z), \quad j = 0, \dots, d-1, \quad (4.19)$$

and the following invariant

$$\prod_k p_j(\omega^k) = \prod_k p_0(\omega^k), \quad j = 0, \dots, d-1. \quad (4.20)$$

Proof: For (4.18) take $m = 0$ in Theorem 4.1. Now re-index to get

$$p_j(\omega^k) = \frac{p_{j+k}(1)\overline{p_k(1)}}{p_0(1)},$$

and take the modulus squared of both sides

$$|p_j(\omega^k)|^2 = \frac{|p_k(1)|^2}{p_0(1)} \frac{|p_{j+k}(1)|^2}{p_0(1)}.$$

But, from $j = 0$ in the first equation,

$$p_0(\omega^k) = \frac{|p_k(1)|^2}{p_0(1)},$$

so that

$$|p_j(\omega^k)|^2 = p_0(\omega^k)p_0(\omega^{j+k}) = p_0(\omega^k)p_0(\omega^j\omega^k),$$

i.e., setting $z = \omega^k$,

$$|p_j(z)|^2 = p_0(z)p_0(\omega^j z).$$

Take the product over k of the re-indexed equation

$$\prod_k p_j(\omega^k) = \frac{1}{(p_0(1))^d} \prod_k p_{j+k}(1)\overline{p_k(1)} = \frac{1}{(p_0(1))^d} \prod_k |p_k(1)|^2,$$

(since each $p_k(1)$ and its conjugate appears exactly *once* in the product). Thus, by (4.19)

$$\prod_k p_j(\omega^k) = \frac{1}{(p_0(1))^d} \prod_{k=0}^{d-1} p_0(1)p_0(\omega^k) = \prod_{k=0}^{d-1} p_0(\omega^k).$$

□

For completeness, we note that the Hermitian condition of Lemma 2.1 can also be succinctly expressed in terms of row symbols.

Lemma 4.2 *Tc is Hermitian if and only if the symbols of c satisfy*

$$\overline{p_j(z)} = p_{-j}(\omega^j z), \quad j = 0, \dots, d-1.$$

Proof: From the definition, we calculate

$$\begin{aligned} \overline{p_j(z)} &= \overline{\sum_r c_{-j,r}(\omega^j z)^{-r}} = \sum_r \overline{c_{-j,r}}(\omega^j z)^r = \sum_k \overline{c_{-j,-k}}\omega^{-jk}z^{-k}, \\ p_{-j}(\omega^j z) &= \sum_r c_{j,r}(\omega^{-j}(\omega^j z))^{-r} = \sum_k c_{jk}z^{-k}, \end{aligned}$$

and so, by equating the coefficients of z^{-k} , the Hermitian condition $c_{jk} = \omega^{-jk}\overline{c_{-j,-k}}$ is equivalent to equality of the above symbols. □

5 The Special Case of $d = 3$

This case already has some interesting geometric features.

Solving the basic equations for the c_{jk} , is also geometrically interesting.

Proposition 5.1 *For $d = 3$, c generates a Heisenberg frame with $v_0 \neq 0$ if and only if*

(a) $p_0(z) = 1 + \overline{c_{01}}z + c_{01}z^2$ and $p_2(z) = \overline{p_1(\omega^2 z)}$ (Hermitian conditions)

(b) $|p_1(z)|^2 = p_0(z)p_0(\omega z)$ (Riesz factorization)

(c) $\prod_{k=0}^2 p_1(\omega^k) = \prod_{k=0}^2 p_0(\omega^k)$ (invariant condition)

(d) $|c_{01}| = |c_{1k}| = \frac{1}{2}$, $k = 0, 1, 2$.

Proof: By Lemma 4.2, the conditions for Tc to be Hermitian are

$$\overline{p_0(z)} = p_0(z), \quad \overline{p_1(z)} = p_2(\omega z), \quad \overline{p_2(z)} = p_1(\omega^2 z).$$

Since $p_0(z) = c_{00} + c_{01}z^2 + c_{02}z$, the first equation is satisfied provided

$$\overline{c_{00}} + \overline{c_{01}}z + \overline{c_{02}}z^2 = c_{00} + c_{02}z + c_{01}z^2 \iff c_{00} \in \mathbb{R}, \quad c_{02} = \overline{c_{01}}.$$

The second and third are equivalent, since substituting ωz for z in the third gives

$$\overline{p_2(\omega z)} = p_1(\omega^2(\omega z)) = p_1(z).$$

Hence (a) is equivalent to Tc being Hermitian with $c_{00} = 1$, and implies $c_{02} = \overline{c_{01}}$.

By Corollary 4.1, (a),(b),(c), (d) hold for a Heisenberg frame with $v_0 \neq 0$. For the converse, suppose that (a),(b),(c), (d) hold. Then by (a), Tc is Hermitian with $c_{00} = 1$, and $c_{02} = \overline{c_{01}}$, so that (d) gives $|c_{jk}| = \frac{1}{2}$, $(j, k) \neq (0, 0)$. Hence Lemma 2.1, gives

$$P_{Tc}(\lambda) = \lambda^3 - \lambda^2 + 0\lambda + a_0, \quad a_0 = -\det(Tc).$$

In view of Theorem 2.1, with condition (iv)', we need only show that $\det(Tc) = 0$.

Since $d(Tc)_{jk} = p_{j-k}(\omega^k)$, condition (a) gives

$$3(Tc) = \begin{bmatrix} p_0(1) & p_2(\omega) & p_1(\omega^2) \\ p_1(1) & p_0(\omega) & p_2(\omega^2) \\ p_2(1) & p_1(\omega) & p_0(\omega^2) \end{bmatrix} = \begin{bmatrix} p_0(1) & \overline{p_1(1)} & \overline{p_1(\omega^2)} \\ \frac{p_1(1)}{\overline{p_1(\omega^2)}} & p_0(\omega) & \overline{p_1(\omega)} \\ p_1(\omega) & p_1(\omega) & p_0(\omega^2) \end{bmatrix}.$$

Since $\overline{p_0(z)} = p_0(z)$, the invariant condition gives $\prod_k p_1(\omega^k) = \prod_k \overline{p_1(\omega^k)} = \prod_k \overline{p_0(\omega^k)}$, and we calculate and so

$$\begin{aligned} 3 \det(Tc) &= p_0(1) \{p_0(\omega)p_0(\omega^2) - |p_1(\omega)|^2\} - \overline{p_1(1)} \{p_1(1)p_0(\omega^2) - \overline{p_1(\omega^2)}p_1(\omega)\} \\ &\quad + p_1(\omega^2) \{p_1(1)p_1(\omega) - \overline{p_1(\omega^2)}p_0(\omega)\} \\ &= 3 \prod_k p_0(\omega^k) - p_0(1)|p_1(\omega)|^2 - p_0(\omega^2)|p_1(1)|^2 - p_0(\omega)|p_1(\omega^2)|^2. \end{aligned}$$

Applying the Riesz factorisation to the last three terms we then obtain

$$3 \det(Tc) = 3 \prod_k p_0(\omega^k) - p_0(1)p_0(\omega)p_0(\omega^2) - p_0(\omega^2)p_0(1)p_0(\omega) - p_0(\omega)p_0(\omega^2)p_0(1) = 0.$$

(need to check $v_0 \neq 0$!) □

We now use Proposition 5.1 to find the solutions for $d = 3$. First we consider the Riesz-type factorisation $|p_1(z)|^2 = p_0(z)p_0(\omega z)$. Note that $1 + \omega + \omega^2 = 0$, and the variable z of our symbols satisfies $z^3 = 1$, $\bar{z} = z^2$. Hence multiplying out gives

$$\begin{aligned} |p_1(z)|^2 &= p_1(z)\overline{p_1(z)} = (c_{20} + c_{21}\omega^2 z^2 + c_{22}\omega z)(\overline{c_{20}} + \overline{c_{21}}\omega z + \overline{c_{22}}\omega^2 z^2) \\ &= (\sum_k |c_{2k}|^2) + (c_{20}\overline{c_{21}} + c_{21}\overline{c_{22}} + c_{22}\overline{c_{20}})\omega z + (c_{20}\overline{c_{22}} + c_{21}\overline{c_{20}} + c_{22}\overline{c_{21}})\omega^2 z^2, \end{aligned}$$

and

$$\begin{aligned} p_0(z)p_0(\omega z) &= (1 + \overline{c_{01}}z + c_{01}z^2)(1 + \overline{c_{01}}\omega z + c_{01}\omega^2 z^2) \\ &= 1 - |c_{01}|^2 + \omega^2(c_{01}^2 - \overline{c_{01}})z + \omega(\overline{c_{01}}^2 - c_{01})z^2. \end{aligned}$$

Hence, equating the coefficients of $1, z, z^2$, gives

$$\begin{aligned} |c_{20}|^2 + |c_{21}|^2 + |c_{22}|^2 &= 1 - |c_{01}|^2, \\ c_{20}\overline{c_{21}} + c_{21}\overline{c_{22}} + c_{22}\overline{c_{20}} &= (c_{01}^2 - \overline{c_{01}})\omega, \\ c_{20}\overline{c_{22}} + c_{21}\overline{c_{20}} + c_{22}\overline{c_{21}} &= (\overline{c_{01}}^2 - c_{01})\omega^2. \end{aligned}$$

Since $|c_{01}|^2 = |c_{1k}|^2 = \frac{1}{4}$, the first equation is automatically satisfied. Further, the second and third are conjugates of each other, and so we have only one equation (for the Riesz-type factorisations)

$$c_{20}\overline{c_{22}} + c_{21}\overline{c_{20}} + c_{22}\overline{c_{21}} = (\overline{c_{01}}^2 - c_{01})\omega^2.$$

Setting $z_{jk} := c_{jk}/|c_{jk}| = 2c_{jk}$, this becomes

$$\frac{1}{4}(z_{20}\overline{z_{22}} + z_{21}\overline{z_{20}} + z_{22}\overline{z_{21}}) = \left(\frac{1}{4}\overline{z_{01}}^2 - \frac{1}{2}z_{01}\right)\omega^2.$$

Since $\overline{z_{01}}^2 = z_{01}^{-2}$, this can be rewritten as

$$z_{20}\overline{z_{22}} + z_{21}\overline{z_{20}} + z_{22}\overline{z_{21}} = \left(\frac{1}{z_{01}^2} - 2z_{01}\right)\omega^2 = \frac{1 - 2z_{01}^3}{z_{01}^2}\omega^2 = \frac{1 - 2(z_{01}/\omega)^3}{(z_{01}/\omega)^2}.$$

Now we set $z := z_{01}/\omega$, so that our equation becomes

$$z_{20}\overline{z_{22}} + z_{21}\overline{z_{20}} + z_{22}\overline{z_{21}} = \frac{1 - 2z^3}{z^2}. \quad (5.21)$$

We proceed to analyze both sides of this equation.

Lemma 5.1 *The curve*

$$\theta \mapsto \frac{1 - 2z^3}{z^2}, \quad z = -e^{i\theta}$$

is a 3-cusped hypocycloid (or deltoid).

Proof: Recall that the standard parametric equations for a hypocycloid with radii a and b with $a > b > 0$ are (see, e.g. [Wik23])

$$\begin{aligned}x(\theta) &= (a - b) \cos(\theta) + b \cos\left(\frac{a - b}{b}\theta\right), \\y(\theta) &= (a - b) \sin(\theta) - b \sin\left(\frac{a - b}{b}\theta\right),\end{aligned}$$

and if $n = a/b$ is an integer, it is n -cusped. Now

$$w := \frac{1 - 2z^3}{z^2} = -2z + z^{-2} = 2e^{i\theta} + e^{-2i\theta},$$

which has Cartesian coordinates

$$\Re(w) = 2 \cos(\theta) + \cos(2\theta), \quad \Im(w) = 2 \sin(\theta) - \sin(2\theta),$$

and so $w(\theta)$ is a 3-cusped hypocycloid with radii $a = 3$ and $b = 1$. □

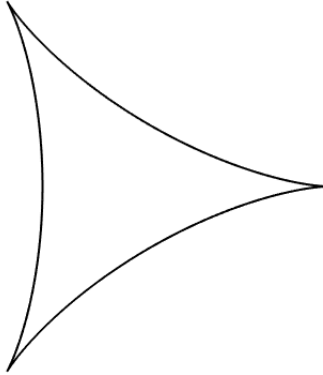


Figure 1: The Hypocycloid

For the left side, note that the product of the three terms

$$z_{20}\overline{z_{22}} \cdot z_{21}\overline{z_{20}} \cdot z_{22}\overline{z_{21}} = 1.$$

Lemma 5.2 *The set of complex numbers*

$$\{z_1 + z_2 + z_3 : z_1 z_2 z_3 = 1, |z_j| = 1\}$$

is the interior and boundary of the 3-cusped hypocycloid given by the right side, i.e.,

$$\theta \mapsto \frac{1 - 2z^3}{z^2}, \quad z = -e^{i\theta}.$$

In particular, points on the boundary have the form

$$z_1 = z_3 = e^{-i\frac{\phi}{2}}, \quad z_2 = e^{i\phi} \quad (\theta = -\frac{\phi}{2}),$$

$$z_1 = z_2 = e^{-i\frac{\phi}{2}}, \quad z_3 = e^{i\phi},$$

or

$$z_2 = z_3 = e^{-i\frac{\phi}{2}}, \quad z_1 = e^{i\phi}.$$

Proof: Since $z_1 z_2 z_3 = 1$, we can write a point w in the set as

$$w = z_1 + z_2 + z_3, \quad z_1 := e^{it}, \quad z_2 := e^{i\phi}, \quad z_3 := e^{-i(t+\phi)}.$$

Now fix ϕ , and let t vary. Let A and B be the points on the hypocycloid for $\theta = -\frac{\phi}{2}$ and $\theta = \pi - \frac{\phi}{2}$, i.e.,

$$A = 2e^{-i\frac{\phi}{2}} + e^{i\phi}, \quad B = -2e^{-i\frac{\phi}{2}} + e^{i\phi}.$$

We claim (cf. Figure 2) that as t varies w traces out the line segment connecting A and B , precisely

$$w = e^{it} + e^{i\phi} + e^{-i(t+\phi)} = \lambda A + (1 - \lambda)B, \quad \lambda = \frac{\cos(t + \frac{\phi}{2}) + 1}{2} \in [0, 1],$$

which we verify by multiplying out

$$\begin{aligned} \lambda A + (1 - \lambda)B &= \lambda(A - B) + B = \frac{e^{i(t+\frac{\phi}{2})} + e^{-i(t+\frac{\phi}{2})} + 2}{4} 4e^{-i\frac{\phi}{2}} - 2e^{-i\frac{\phi}{2}} + e^{i\phi} \\ &= (e^{it} + e^{-i(t+\phi)} + 2e^{-i\frac{\phi}{2}}) - 2e^{-i\frac{\phi}{2}} + e^{i\phi} = e^{it} + e^{i\phi} + e^{-i(t+\phi)}. \end{aligned}$$

Further we note that this line segment is tangent to the point where $\theta = \phi$, i.e.,

$$C = 2e^{i\phi} + e^{-2i\phi}.$$

Indeed the tangent to the hypocycloid at this point is

$$\frac{d}{d\theta}(2e^{i\theta} + e^{-2i\theta})|_{\theta=\phi} = 2e^{i\phi} - 2e^{-2i\phi} = i \sin(\frac{3}{2}\phi)(A - B),$$

which is collinear with the line segment, except when $\phi = 0, \pm\frac{2}{3}\pi$, the three cusps of the hypocycloid.

At the cusp corresponding to $\phi = 0$, $A = 3$ and $B = -1$, and thus the line segment connecting A and B is also “tangent” at that cusp. The other cusps are handled similarly. \square

Thus, from equation (5.21), it follows that $z_{20}\overline{z_{22}}$, $z_{21}\overline{z_{20}}$ and $z_{22}\overline{z_{21}}$ are *boundary* points of the hypocycloid. Solutions may be obtained as follows. Pick one of the boundary solutions, e.g., where $z_1 = z_3$, so that

$$z_{20}\overline{z_{22}} = e^{-i\frac{\phi}{2}}, \quad z_{21}\overline{z_{20}} = e^{i\phi}, \quad z_{22}\overline{z_{21}} = e^{-i\frac{\phi}{2}}, \quad z = -e^{-i\frac{\phi}{2}}.$$

These can be solved using one of them as a free parameter, i.e.,

$$z_{22} = e^{i\frac{\phi}{2}} z_{20} \text{ and } z_{21} = e^{i\phi} z_{20}, \quad |z_{20}| = 1.$$

In this way we arrive at a continuum of parameterized solutions for the overlaps of a SIC in dimension $d = 3$.

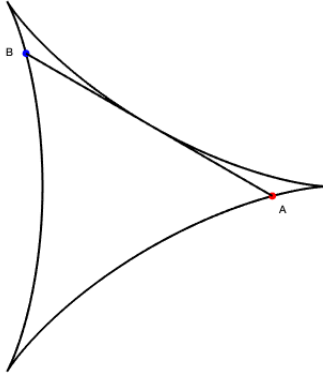


Figure 2: Points A and B on the Hypocycloid

6 Closing Comment

It has sometimes been remarked that the overlaps are zeros of a self-reciprocal polynomial ($z^n p(1/z) = p(z)$) with integer coefficients. The fact that the coefficients are integers is notable and perhaps important. However being self-reciprocal is *not*. Indeed if they are roots of a polynomial $p(z)$ of degree n , then they are also automatically roots of $q(z) := p(z) \times z^n p(1/z)$ and this latter polynomial is self-reciprocal.

7 Acknowledgement

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